

Basic Research of Tritium Confinement

D J Campbell

ITER Organization, Cadarache

Acknowledgements:

Many colleagues in the JET and TFTR teams



Summary of Presentation

- **The aim and planning of DT experiments in present tokamaks**
- **Major results of DT experiments in JET and TFTR:**
 - Plasma energy confinement
 - Particle transport studies
 - MHD stability
 - α -particle behaviour
 - Fusion power production
- **Conclusions**

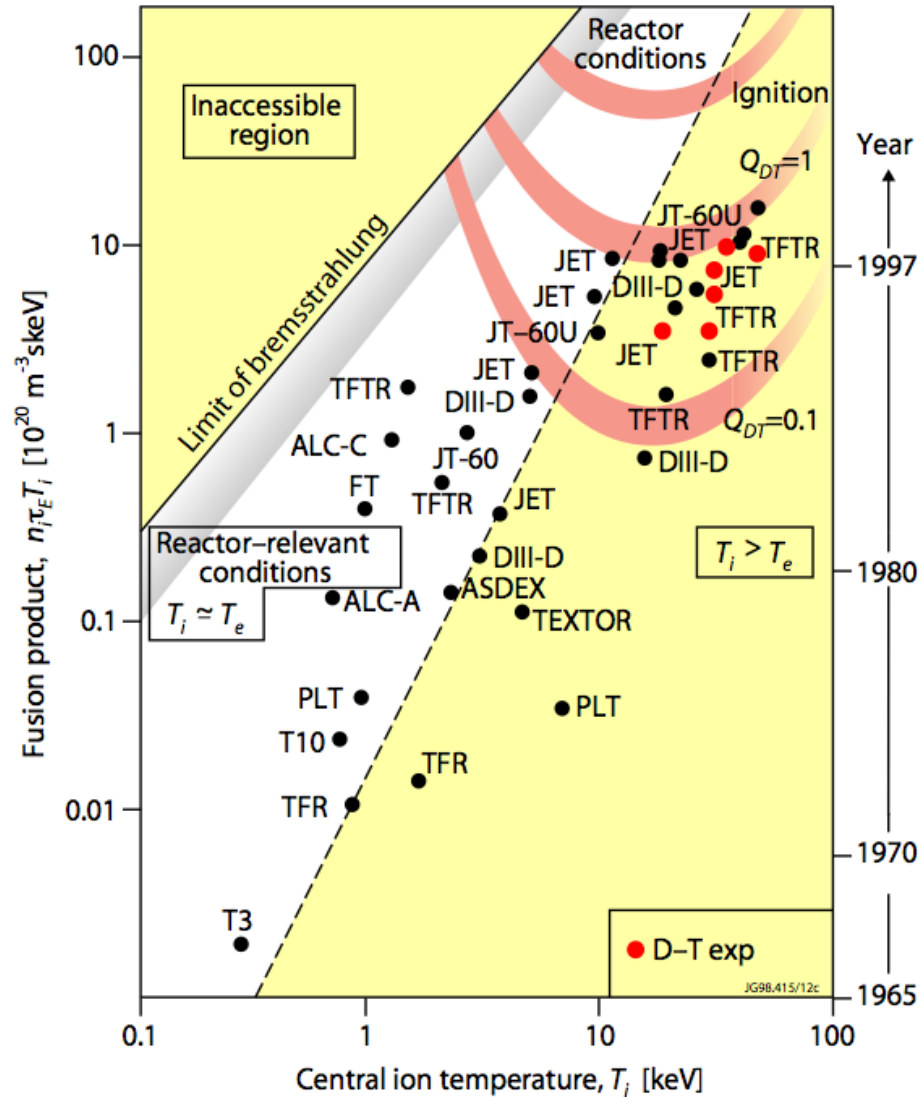
Aim of DT Experiments

- Over the last 4 decades, fusion research has made significant progress towards the goal of energy production:

$$Q = \frac{\text{Fusion Power}}{\text{Input Power}} \sim n_i T_i \tau_E$$

- DT experiments allow:**

- direct test of predictions of fusion performance
- analysis of behaviour of tritium in the tokamak environment
- studies of plasma behaviour with DT mixtures
- investigation of α -particle behaviour and impact on plasmas





Planning DT Experiments

Progress from DD to DT experiments is a major (and exciting!) step for the magnetic fusion programme

- **DT fuel brings a new approach to the organization of the tokamak experimental programme:**
 - Tritium is itself radioactive
 - Limited amounts of tritium are stored on-site to limit licensing requirements
 - Amount of tritium trapped inside vacuum vessel must be limited
 - DT fusion reactivity factor of >100 greater than DD reactivity
 - 14 MeV neutrons vs 2.4 MeV neutrons \Rightarrow additional activation products
- \Rightarrow **experimental programme must be planned with great care to minimize use of tritium and activation of the device structure**
- \Rightarrow **rehearsal of plasma scenarios in deuterium and careful development to optimize use of tritium**



Planning DT Experiments

Sequence of DT Programmes

- **A series of DT experiments has been carried out within the tokamak programme:**

- 1991: JET Preliminary Tritium Experiment (PTE) - $n_T/(n_D+n_T) < 11\%$
- 1994-1997: TFTR DT experiments - $n_T/(n_D+n_T) < 90\%$
- 1997: JET Deuterium-Tritium Experiment (DTE) - $n_T/(n_D+n_T) < 90\%$
- 2003: JET Trace Tritium Experiment (TTE) - $n_T/(n_D+n_T) < 3\%$

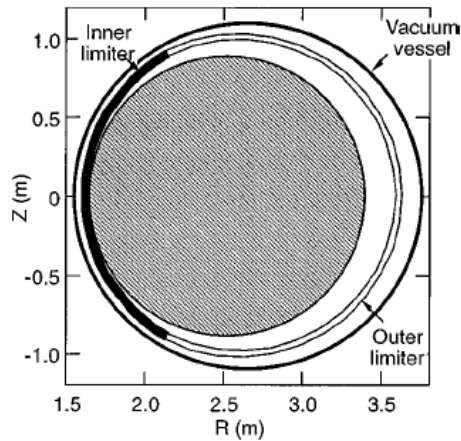
⇒ **provided gradually increasing experience of handling tritium in quantities relevant to tokamak experimental requirements**

⇒ **allowed wide range of physics studies to be developed in DT plasmas in a range of plasma regimes**

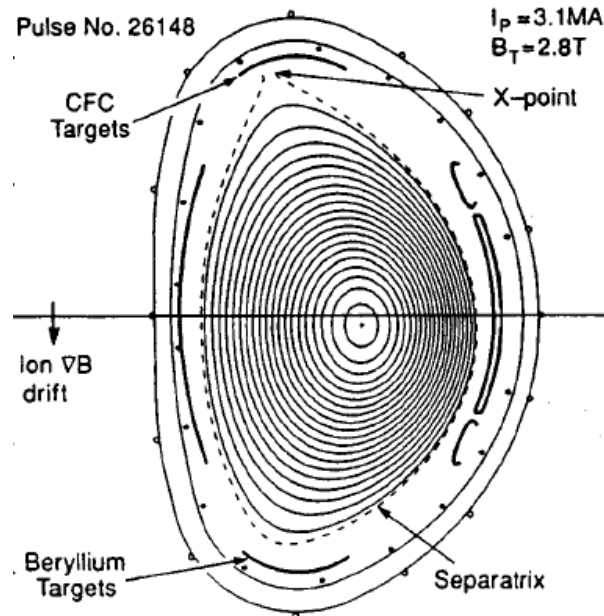
⇒ **produced maximum fusion power of 16MW**

Planning DT Experiments

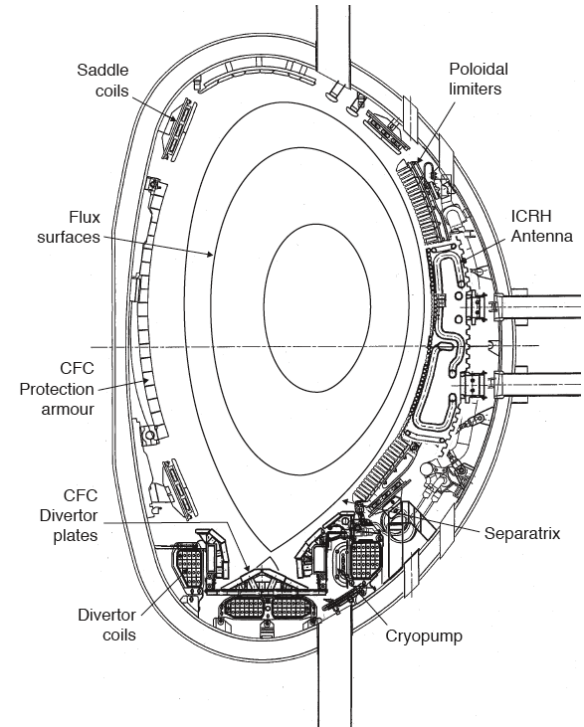
Experimental Configurations



TFTR 1994-1997



JET 1991



JET 1997, 2003

- **TFTR and JET use different magnetic configurations:**

- TFTR DT experiments in limiter plasmas: L-mode, “supershot”, ITB scenarios
- JET DT experiments in diverted plasmas: L-mode, H-mode, ITB scenarios



Planning DT Experiments

Fuelling DT Plasmas

- **Optimization of the fusion power in a DT plasma raises the requirement to be able to control the ratio D:T (ideally 50:50) - implies a need for:**
 - measuring the “fuel ratio”
 - providing controlled amounts of deuterium and tritium
- **Generally the plasmas which will be discussed here were fuelled by a mixture of**
 - gas fuelling
 - neutral beam injection
 - recycling from plasma facing surfaces
- **Control of the fuel content of the plasma facing surfaces was therefore an important aspect of the DT experiments carried out in JET and TFTR**



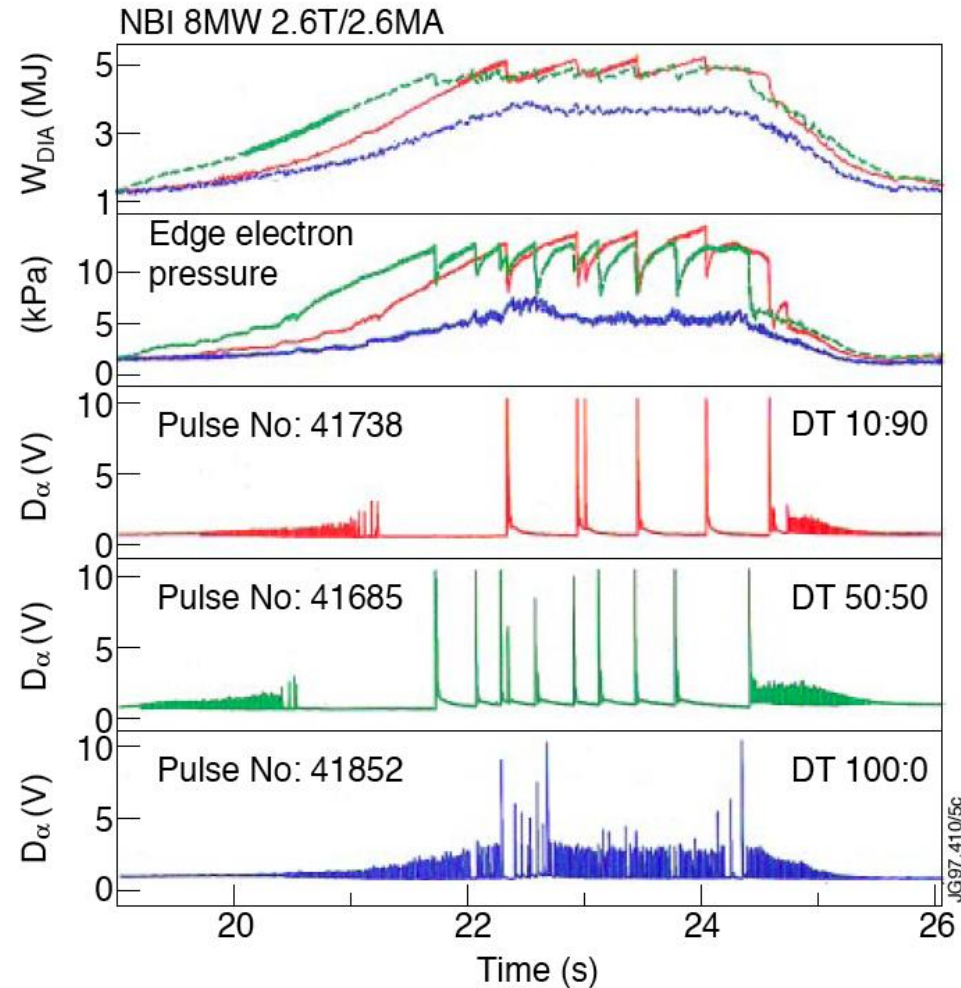
Experimental Results in DT Plasmas: Plasma Energy Confinement

Plasma Energy Confinement

Tritium in Plasmas

- Presence of tritium in plasmas has an effect on plasma behaviour which is independent of the production of fusion power

⇒ **Isotope effect**



JET H-modes



Plasma Energy Confinement

1. TFTR “Supershots”

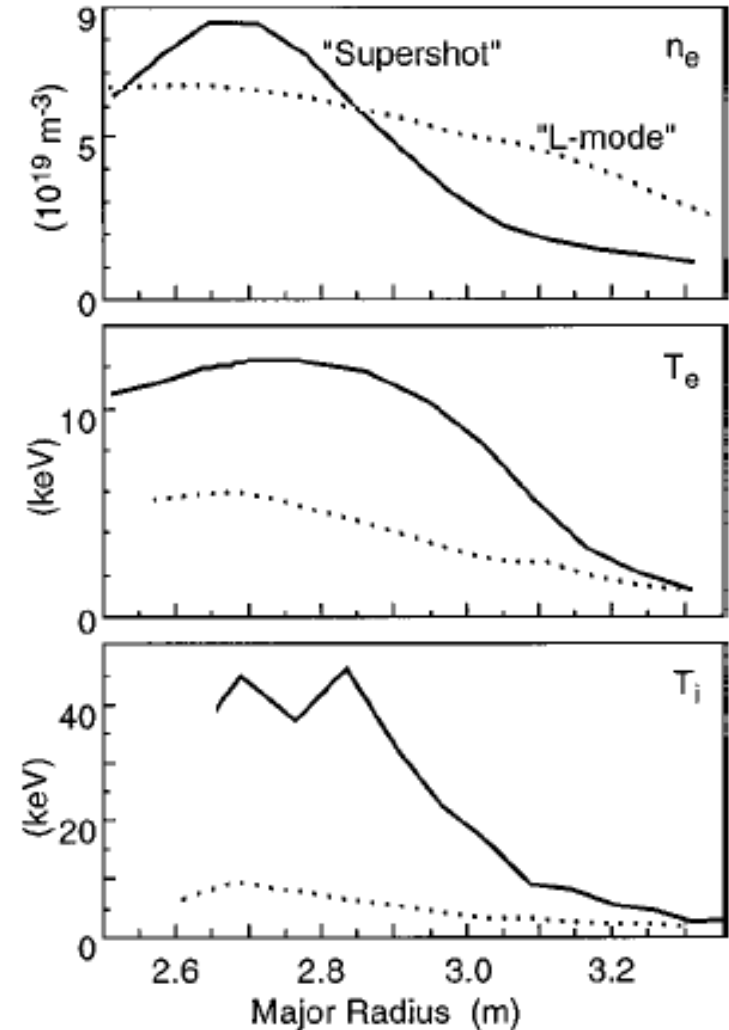
- “Supershots” are limiter plasmas in TFTR with well conditioned walls and high power neutral beam injection:
 - large reduction of central transport with peaked density and temperature profiles
 - associated with rotational shear suppression of turbulence in plasma core

- DT experiments showed:

$$\tau_E^{\text{thermal}} \propto \langle A \rangle^{0.89}$$

$$\chi_i^{\text{tot}} \propto \langle A \rangle^{-2.6}$$

- ie a strong dependence on the isotope mass, A



J D Strachan et al, Phys Rev Lett **58** 1004 (1987)



Plasma Energy Confinement

2. JET L-modes

- **JET L-mode experiments :**

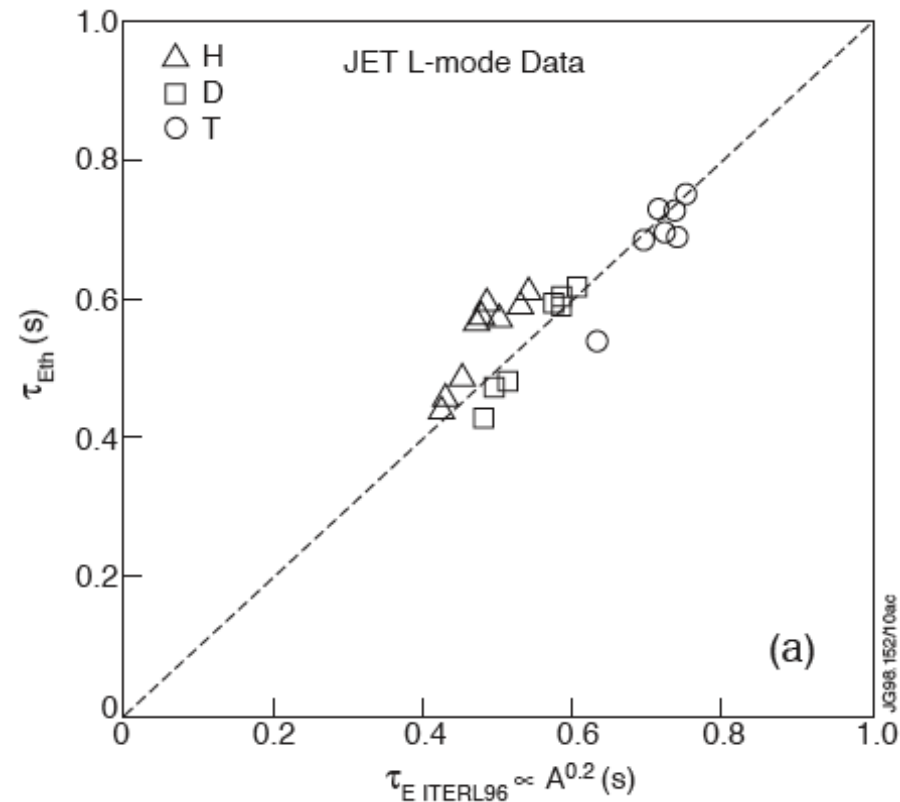
$$\tau_E^L \propto \langle A \rangle^{0.0} - \langle A \rangle^{0.2}$$

- comparable to L-mode global confinement scaling

- **For comparison, TFTR L-mode plasmas yielded:**

$$\tau_E^L \propto \langle A \rangle^{0.3} - \langle A \rangle^{0.5}$$

- A possible explanation lies in low input power of JET pulses - possible source of the discrepancy as plasmas are “quasi-ohmic”.



J G Cordey et al, Nucl Fusion **39** 301 (1999)



Plasma Energy Confinement

3. JET H-modes

- Simple analysis of JET ELMy H-mode experiments suggested:

$$\tau_E^H \propto \langle A \rangle^{0.16}$$

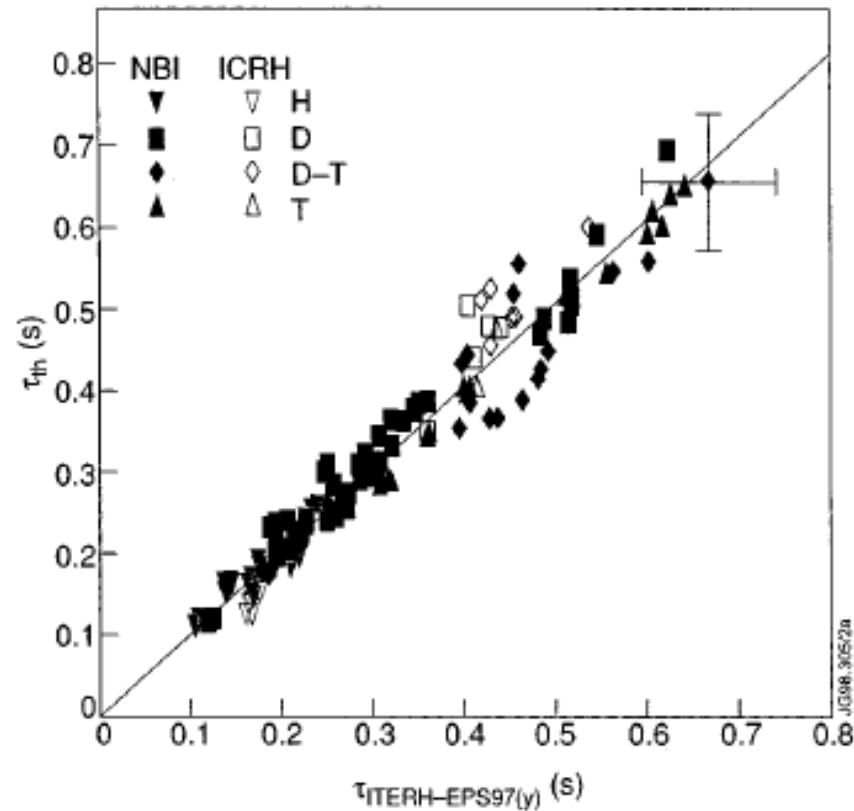
- However, better matching of power and density data yielded:

$$\tau_E^H \propto \langle A \rangle^{0.03}$$

- Finally, when **core** and **pedestal** energy contributions separated:

$$W_{\text{core}}^H \propto \langle A \rangle^{-0.17}$$

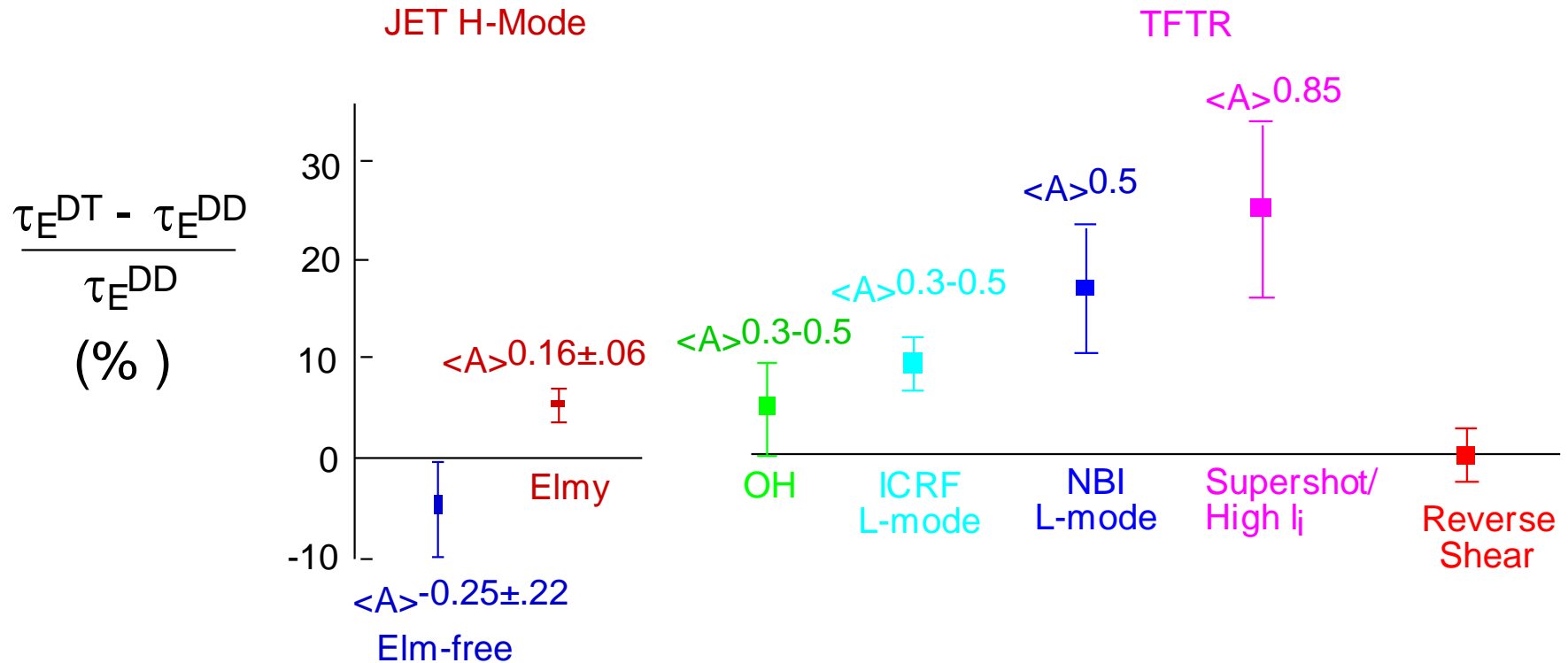
- this is close to $A^{-0.2}$ scaling expected from gyro-Bohm transport model



J G Cordey et al, Nucl Fusion **39** 301 (1999)



Plasma Energy Confinement Overview



- Overall, the isotope dependence of confinement has been found to vary widely, depending on plasma operating regime:
 - indicates different processes influencing confinement and their varying importance in different plasma regimes



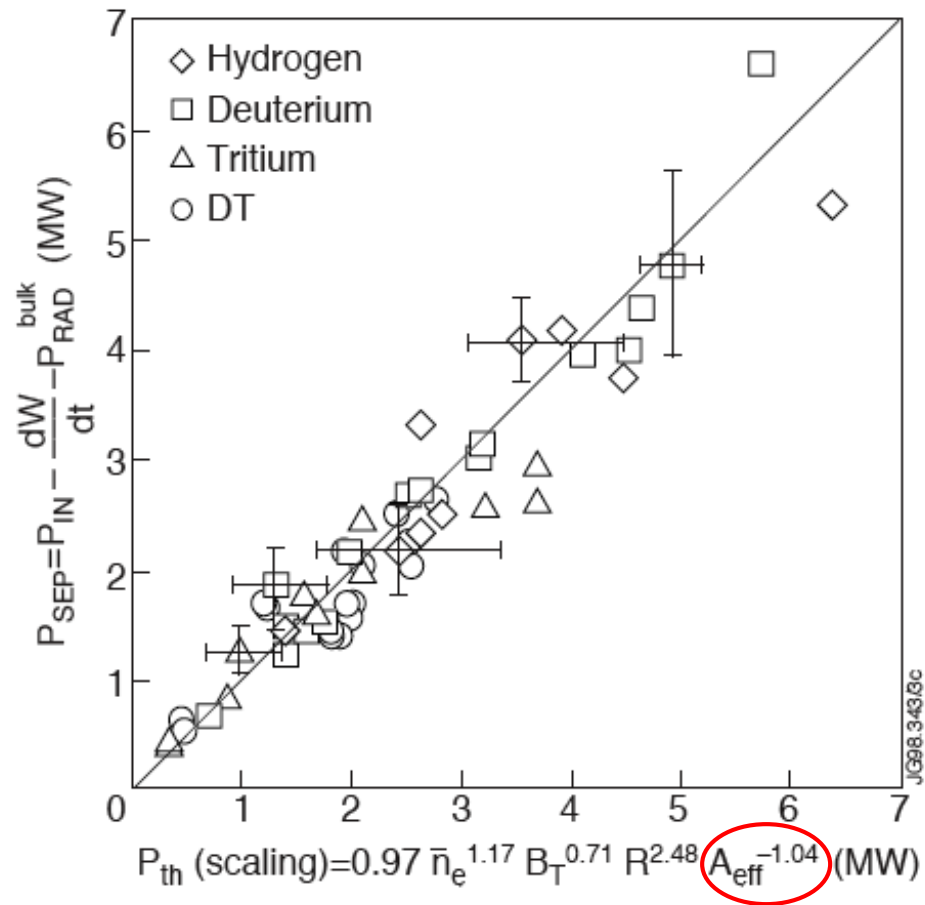
Plasma Energy Confinement

H-mode Power Threshold

- JET analysis of the power required to access the H-mode confirmed that:

$$P_{th} \propto \langle A \rangle^{-1}$$

- This result is important for ITER in that it indicates that access to the H-mode will be easiest in DT operation



E Righi et al, Nucl Fusion **39** 309 (1999)

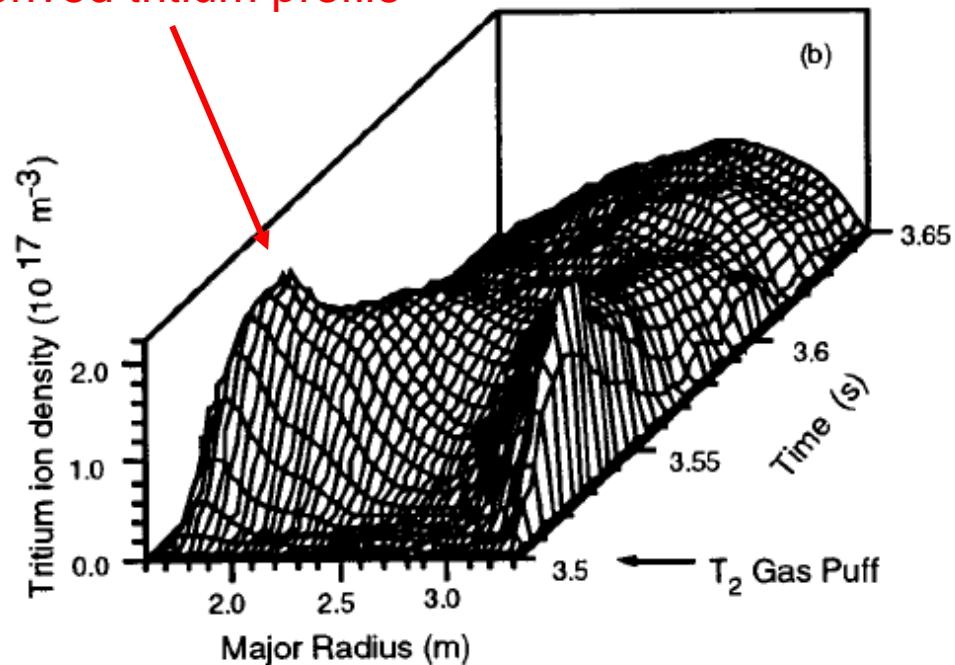
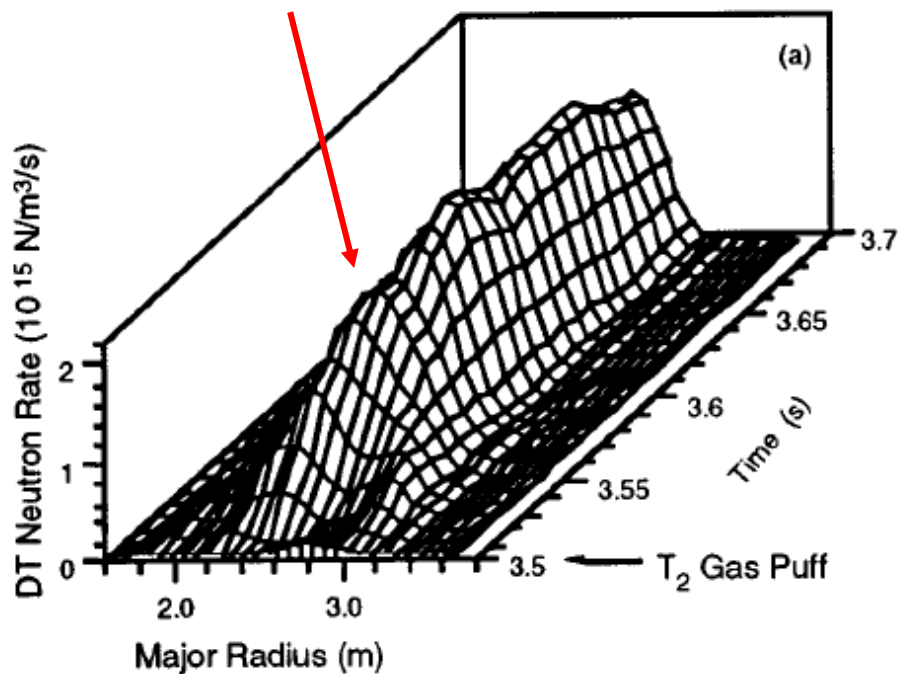


Experimental Results in DT Plasmas: Particle Transport Studies

Particle Transport Studies Principle

14MeV neutron signal

Derived tritium profile

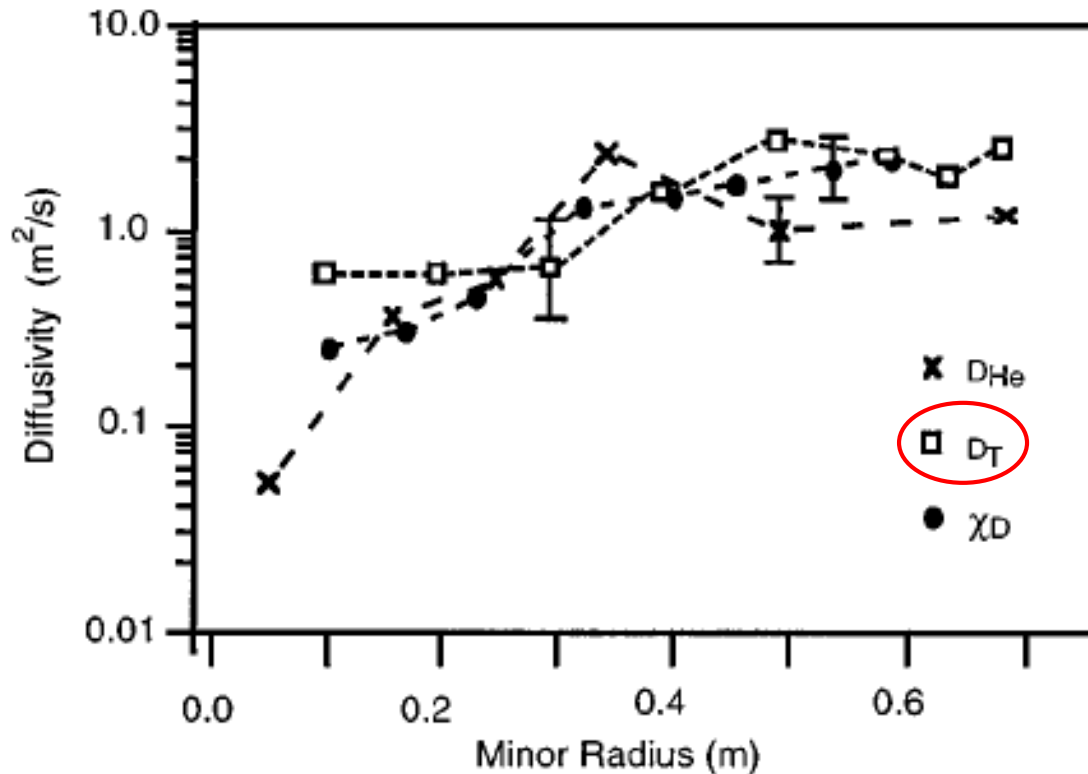


P C Efthimion et al, Phys Rev Lett **75** 85 (1995)

- **Puffing a small amount of tritium into the edge of a deuterium plasma allows hydrogenic particle transport to be investigated:**
 - The 14MeV neutron signal provides a distinct marker for the location of the tritium as it diffuses into the plasma

Particle Transport Studies

1. TFTR “Supershots”



P C Efthimion et al, Phys Rev Lett **75** 85 (1995)

- In “supershot” plasmas, the tritium particle diffusivity was very similar to the electron and helium particle diffusivity:

Particle Transport Studies

2. JET non-dimensional scans

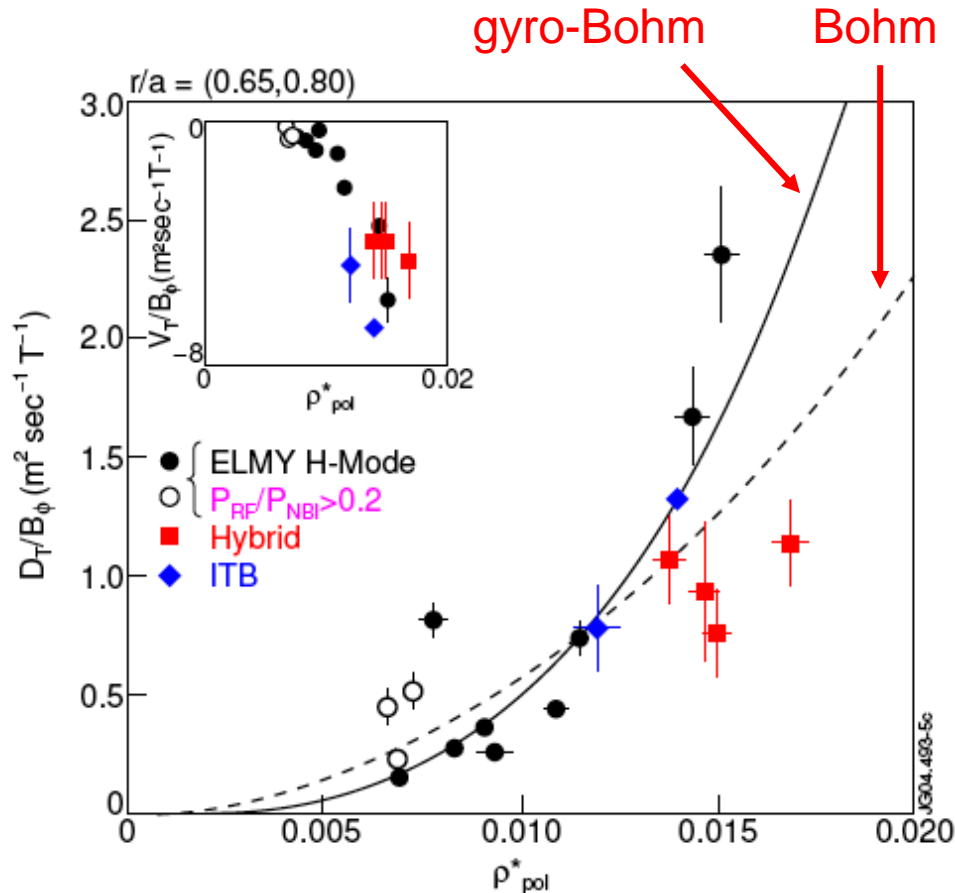
- Studies of tritium particle diffusivity in “non-dimensional” studies showed that in **ELMy H-modes**:

- particle diffusivity in the radial region $0.65 < r/a < 0.8$ was well correlated with gyro-Bohm scaling:

$$\frac{D_T}{B_\phi} \propto \rho_{pol}^{*3}$$

- rather than Bohm scaling:

$$\frac{D_T}{B_\phi} \propto \rho_{pol}^{*2}$$



D Stork et al, IAEA-20, Vilamoura (2004)



Experimental Results in DT Plasmas: MHD Stability



Alfvén Eigenmodes

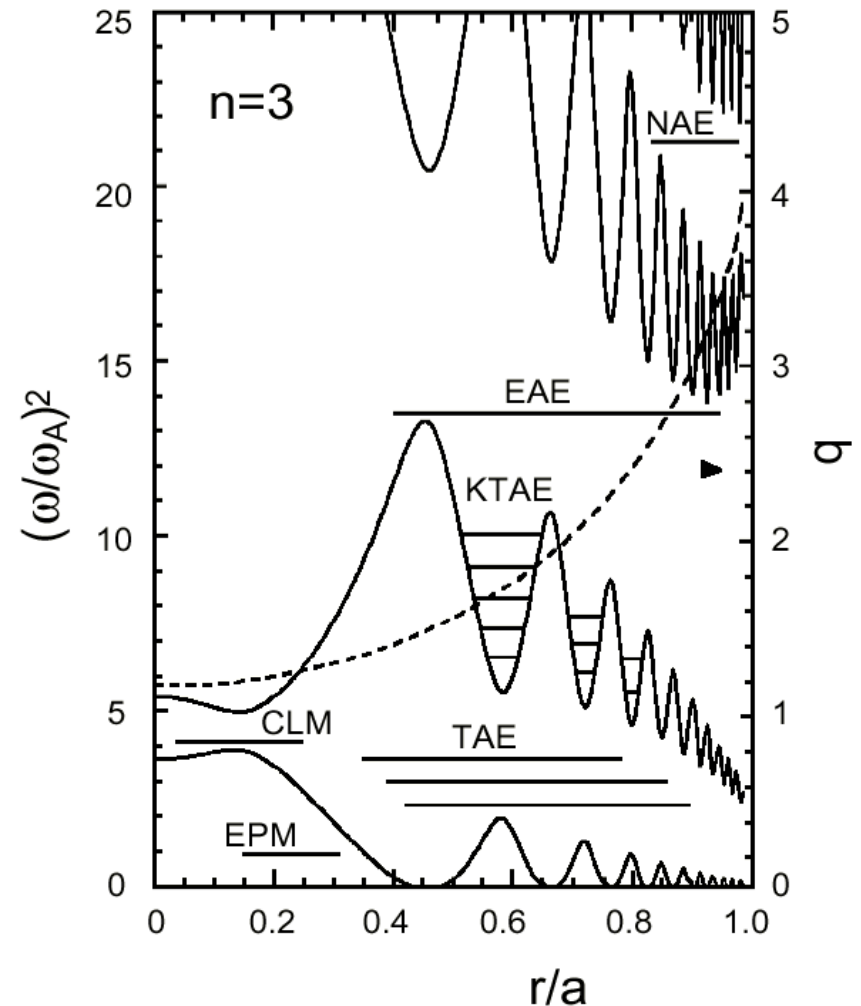
- In a tokamak plasma, the Alfvén wave continuum splits into a series of bands, with the gaps associated with various features of the equilibrium:
 - a series of discrete frequency Alfvén eigenmodes can exist in these gaps:
 - toroidicity-induced (TAE) gap created by toroidicity
 - ellipticity-induced (EAE) gap created by elongation
 - triangularity-induced (NAE) gap created by additional non-circular effects
 - beta-induced (BAE) gap created by field compressibility
 - kinetic toroidal (KTAE) gap created by non-ideal effects such as finite Larmor radius
- ... and others!
- These modes can be driven unstable by the free energy arising from energetic particle populations with velocities above the Alfvén velocity, eg α -particles

Alfvén Eigenmode Spectra

- Characteristic **Alfvén frequency** defined:

$$\omega_A = \frac{v_A}{qR}$$

- TAE centred near $\omega_A/2$
- EAE centred near ω_A
- NAE centred near $3\omega_A/2$
- BAE centred near $\omega_{TAE}/2$
- Each AE is specified by a specific **toroidal mode number, n**:
 - for each n, mode structure consists of a superposition of poloidal harmonics, with the m-harmonic peaking near $q=m/n$



G J Kramer et al, Phys Rev Lett **80** 2594 (1998)



Key Energetic Ion Parameters

The influence of energetic ion populations on plasma stability can be expressed through a small number of parameters

- $\frac{\delta_f}{a} = \frac{q}{\varepsilon^{0.5}} \frac{r_f}{a}$ - normalized half-width of fast ion banana orbit
- n_f / n_e - fractional density of fast ions
- $\beta_f(0)$ - normalized axial fast ion pressure
- $\max |R \cdot \nabla \beta_f|$ - dimensionless fast ion pressure gradient
- $v_f / v_A(0)$ - ratio of fast ion velocity to central Alfvén velocity

$$v_A = \frac{B}{\sqrt{\mu_0 \rho_{\text{mass}}}}$$



α -Particle Physics in Advanced Scenarios

- Excitation of AEs and their influence on α -particle confinement is a central question for viability of advanced scenarios
- α -particle parameters in TFTR and JET might allow access to relevant range where **α -driven instabilities** and their influence on **α -particle transport** could be studied:

Parameter	α 's (TFTR)	α 's (JET)	α 's (ITER)
$P_f(0)$ [MWm ⁻³]	0.3	0.16	0.44
δ/a	0.3	0.34	0.08
$n_f(0)/n_e(0)$ [%]	0.3	0.17	0.8
$\beta_f(0)$ [%]	0.26	0.3	1.1
$\langle\beta_f\rangle$ [%]	0.03	0.04	0.16
$\max R \cdot \nabla \beta_f $ [%]	2	1.6	8
$v_f / v_A(0)$	1.6	1.4	1.8

- Higher power density in reactor could be even more challenging
⇒ ITER programme aims to move in this direction



Alfvén Eigenmode Stability

Linear stability:

- **Pressure gradient of resonant particles with $v \sim v_A$ provides source of free energy which excites the mode:**
 - both passing and trapped particles can resonate with the AE
 - a resonant sideband also exists at $v_A/3$ for TAEs
 - several damping mechanisms exist which complicate estimation of instability thresholds
 - experiments in “advanced scenarios” with **non-monotonic q-profiles** show a rich population of AEs can be excited

Non-linear behaviour:

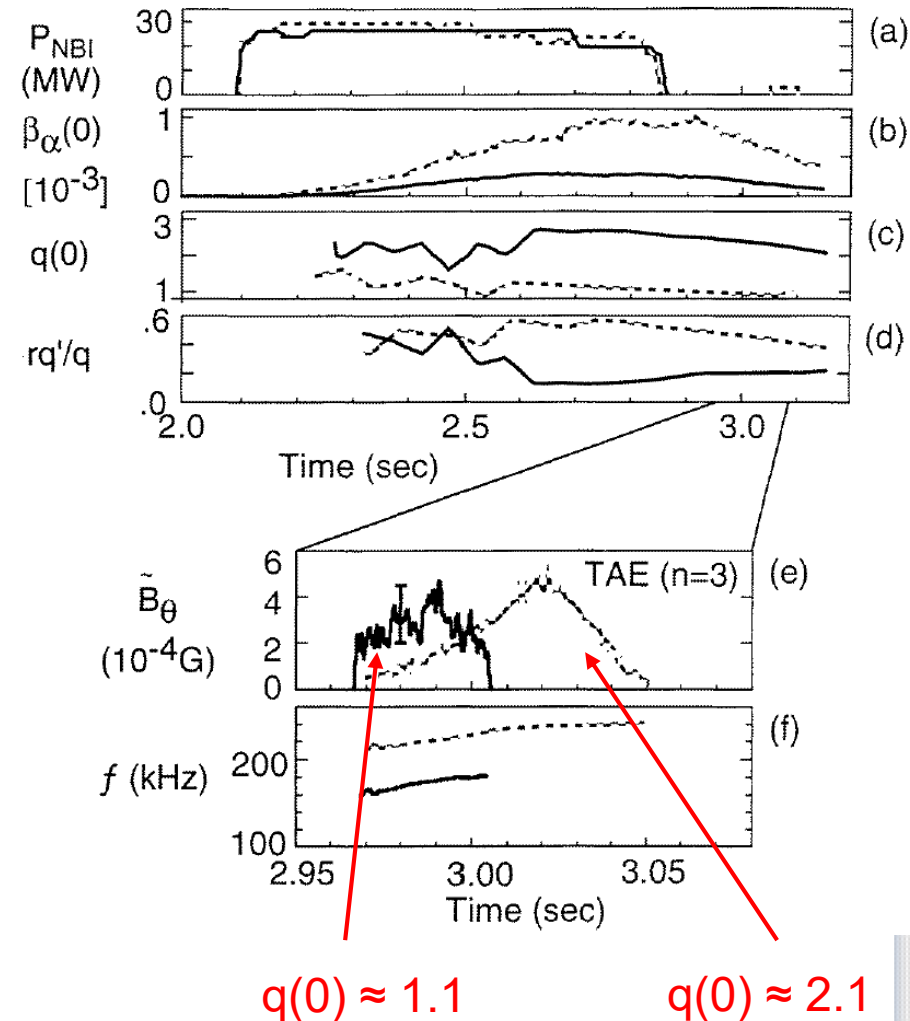
- **Redistribution of resonant particles can occur by finite amplitude waves**
- **Overlap of multiple modes can lead to enhancement of energetic ion transport**



Alfvén Eigenmode Observation

- In both TFTR and JET DT experiments, excitation of AEs by α -particles proved elusive. TFTR succeeded by:
 - raising central q , to increase instability drive
 - observing in post-NBI phase, where (beam) ion Landau damping is reduced, but α -drive remains
 - TAE may have been generated through excitation of so-called “Cascade” mode
- JET reported observations of ICRH-driven TAEs for which α -particles contributed to instability drive

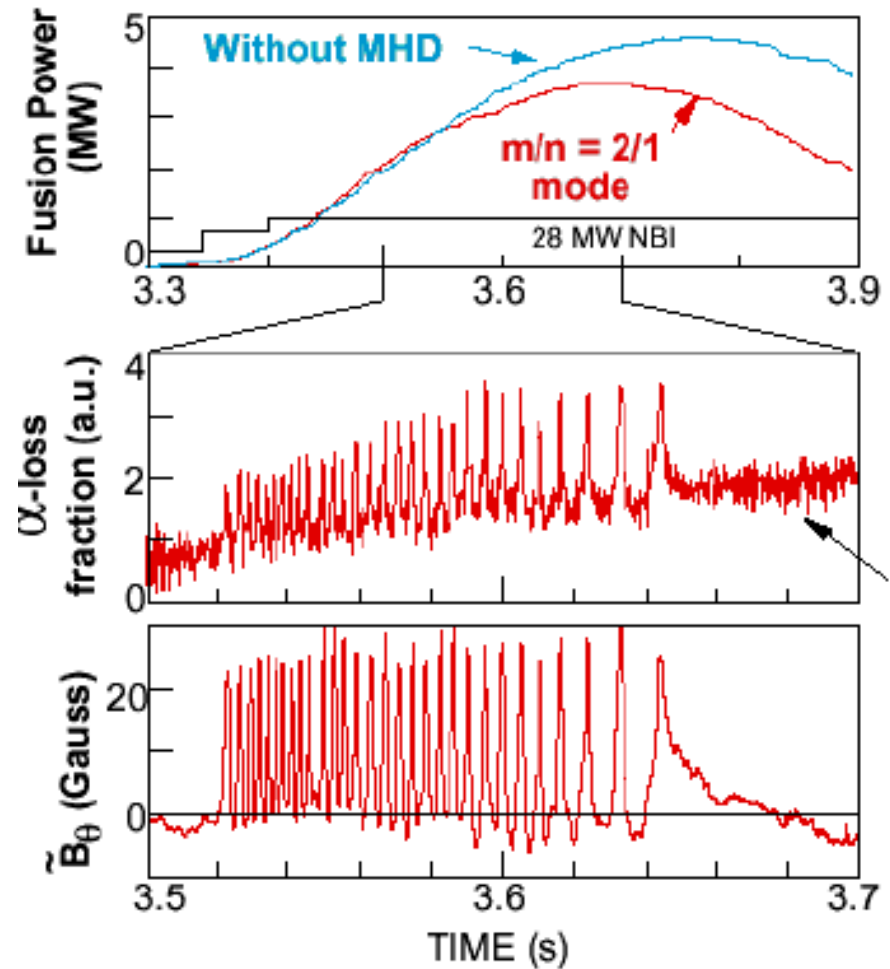
R Nazikian et al, Phys Rev Lett **78** 2976 (1997)



Alpha Particles and MHD

Enhanced losses through instabilities

- In TFTR, strong anisotropic loss was observed as the NTM mode was rotating:
 - Disruptions, kinetic ballooning modes and sawteeth were also observed to generate α -particle losses
- Typically, sawteeth produced a large radial redistribution of α -particles



K M McGuire et al, IAEA-16, Montreal (1996)

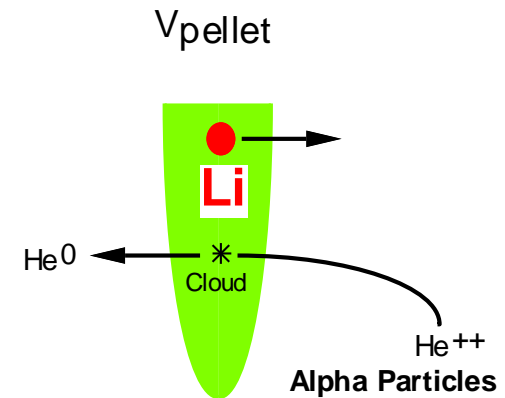
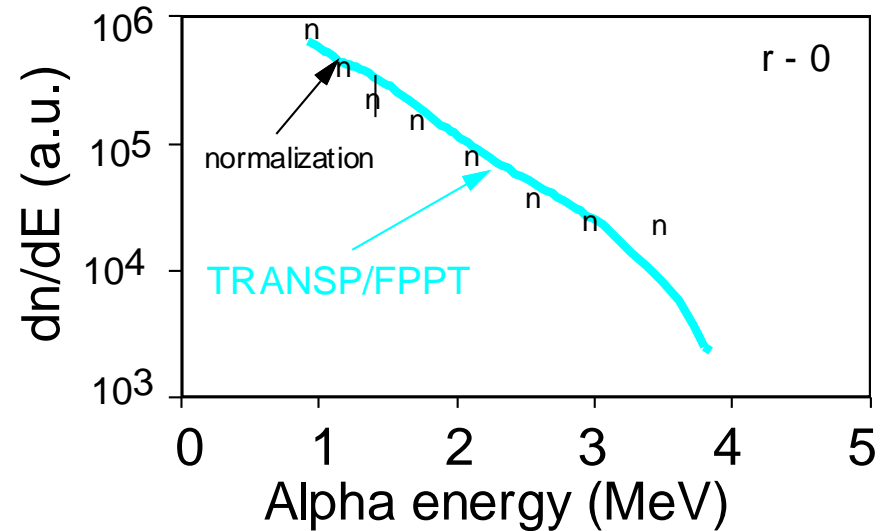


Experimental Results in DT Plasmas: α -Particle Behaviour

α -Particle Behaviour

Classical slowing down

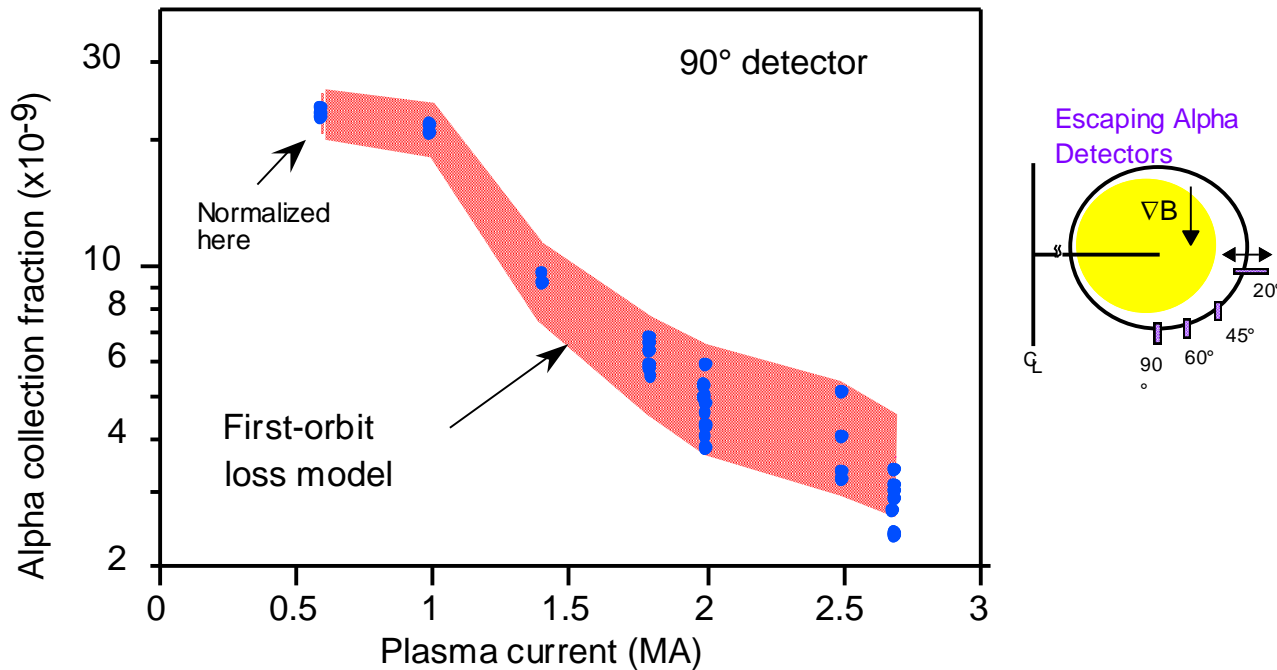
- In TFTR, measurements were made of the slowing down spectrum of α -particles:
 - Both charge exchange spectroscopy on lithium pellets and charge exchange spectroscopy on neutral beams confirmed classically calculated slowing down spectrum
- Confirms interaction between 3.5MeV α -particles and thermal plasma is as expected theoretically



R K Fischer et al, Phys Rev Lett **78** 846 (1995)

α -Particle Behaviour

Classical confinement

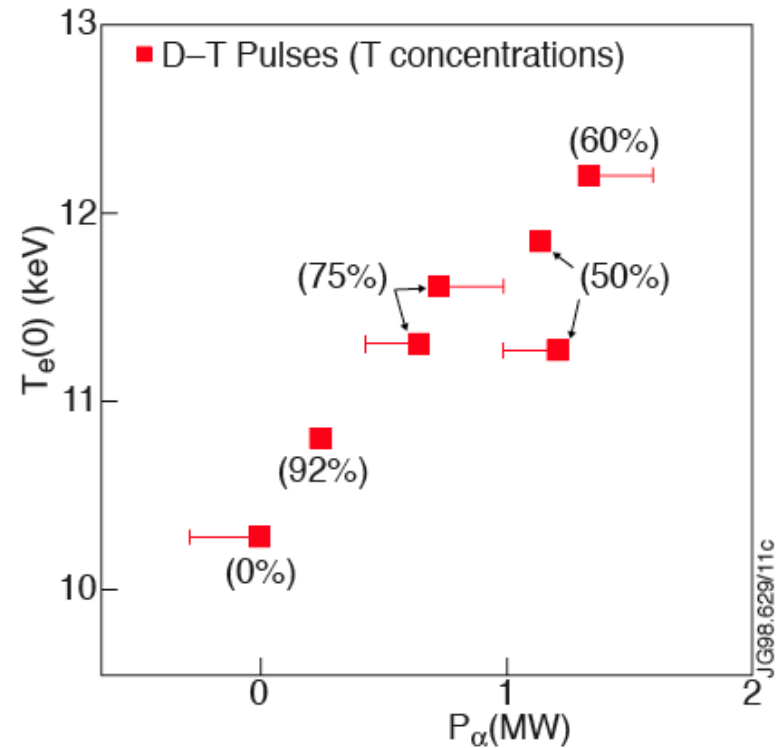
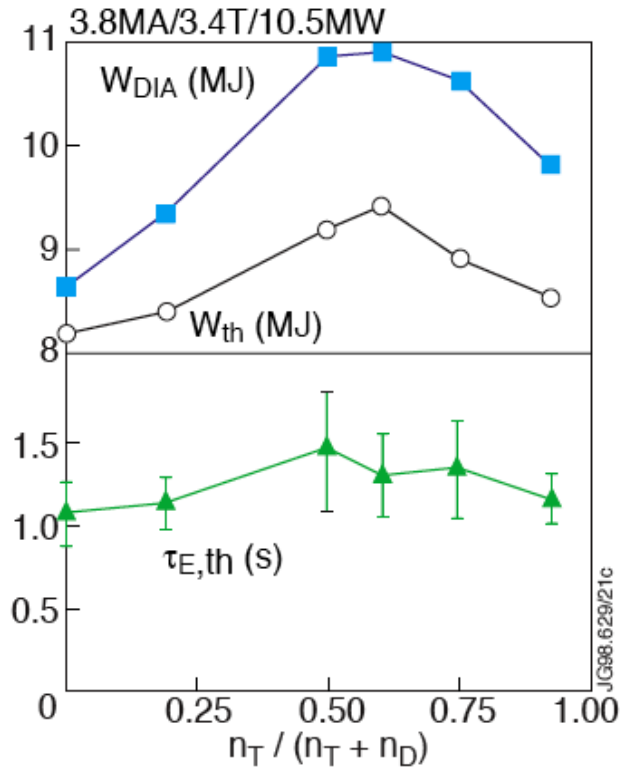


- **Measurements made of escaping α -particles at 90° to midplane as a function of plasma current:**
 - escaping α -particle flux was observed to be consistent with classical first orbit losses

R J Hawryluk et al, IAEA-15, Seville (1994)

α -Particle Behaviour

Electron heating



P R Thomas et al, Phys Rev Lett **80** 5548 (1998)

- In JET, a controlled scan of tritium concentration demonstrated that maximum stored energy and electron temperature correlated with maximum fusion power and optimum DT mix:
 - confirms that α -particles heat the electrons efficiently as expected



Experimental Results in DT Plasmas: Fusion Power Production



Fusion Power Production

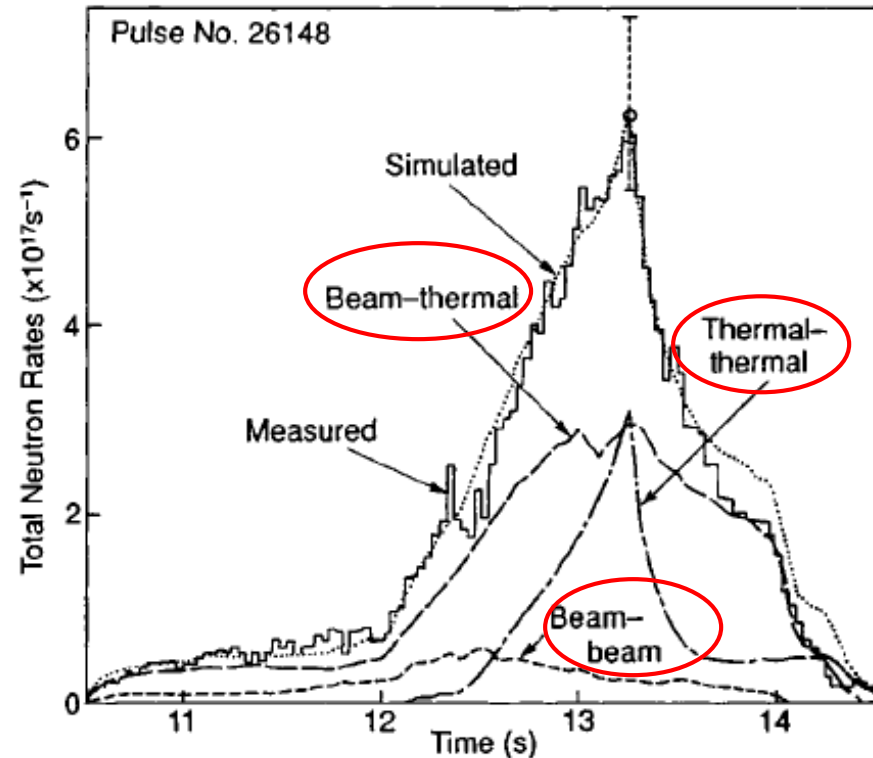
- **During the 1990s a series of experiments were designed to optimize DT fusion power production from tokamak plasmas:**
 - To achieve highest possible fusion power and highest Q values
 - To investigate highest powers which could be achieved in stationary operation
 - To allow α -particle heating of plasmas to be detected
- **Plasmas were optimized to achieve the maximum total power achievable:**
 - typically this involved developing low density target plasmas into which high power neutral beams were injected:
 - JET: hot-ion H-mode
 - TFTR: supershot



Fusion Power Production Contributions to Fusion Reactivity

- In present devices, the fusion power production is made up of several components:
 - reactions between thermal plasma particles (thermal-thermal)
 - reactions between injected beam ions and thermal particles (beam-thermal)
 - reactions between injected beam ions themselves (beam-beam)
- Figure shows a comparison between measured total neutron yield and calculated components of neutron production in JET PTE (11% tritium)

JET Team, Nucl Fusion **32** 187 (1992)

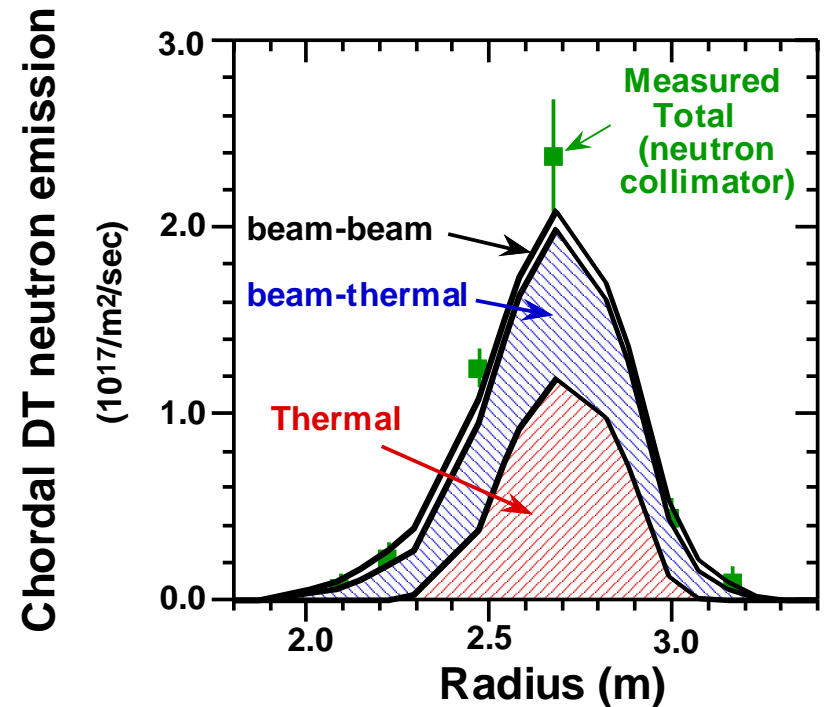


JET PTE (1991)
1.7MW fusion power



Fusion Power Production Contributions to Fusion Reactivity

- Analysis of TFTR supershots shows a similar behaviour:
 - for a measurement chord through the plasma core, thermal and beam-thermal reactions contribute ~50% of total
 - towards plasma edge, lower plasma temperatures imply that beam-thermal contribution gradually dominates - but at significantly lower reaction rate



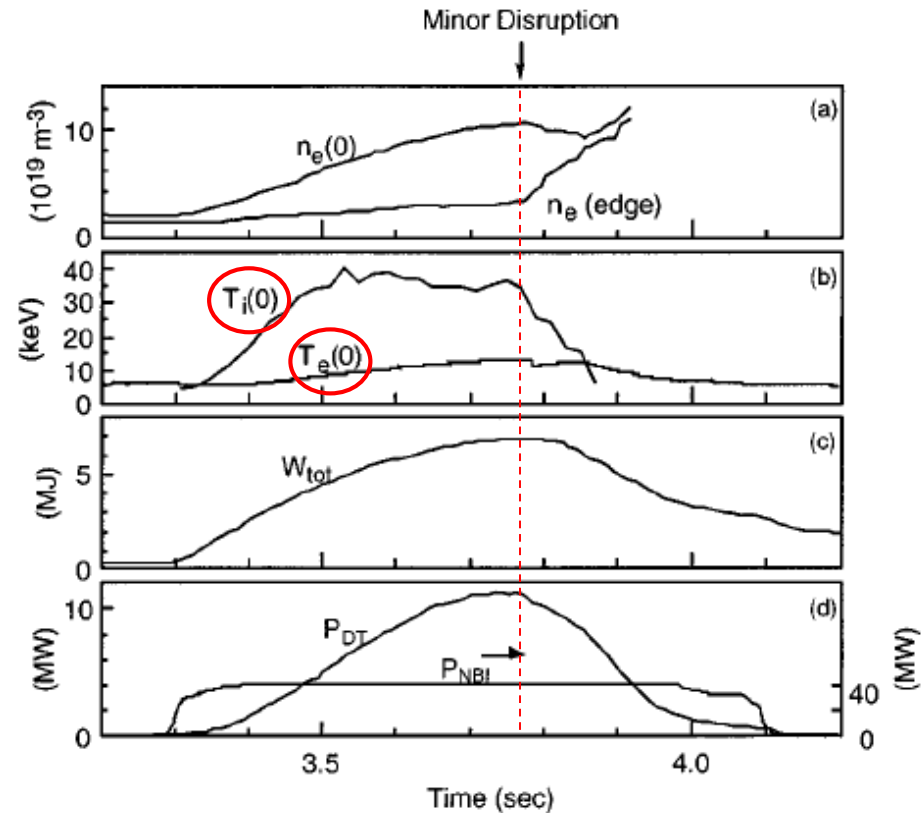
TFTR DT (1994-1997)

R V Budny et al, Nucl Fusion **35** 1497 (1995)

Fusion Power Production

TFTR Supershot

- The best TFTR supershot produced 11MW of DT fusion power with ~40MW of NB input power:
 - $Q = 0.27$
- The high fusion power pulse lasts less than 1s due to a combination of MHD instability and a subsequent influx of impurities (“carbon bloom”)

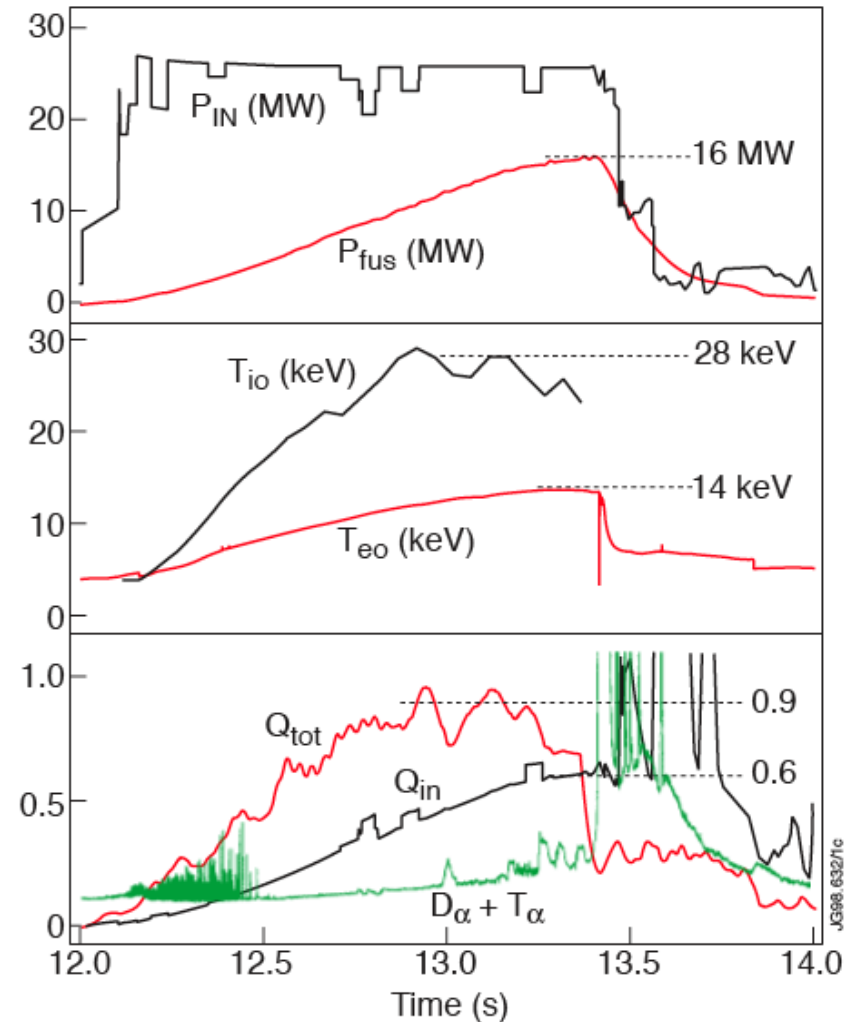


R J Hawryluk et al, Phys Plasmas **5** 1577 (1998)

Fusion Power Production

JET Hot-Ion H-mode

- JET hot-ion H-modes follow a similar time behaviour to TFTR supershots:
 - $P_{\text{fus}} = 16\text{MW}$
 - $Q = 0.6$
- The high fusion power pulse again lasts about 1s due to a combination of MHD instability and a subsequent influx of impurities (“carbon bloom”)

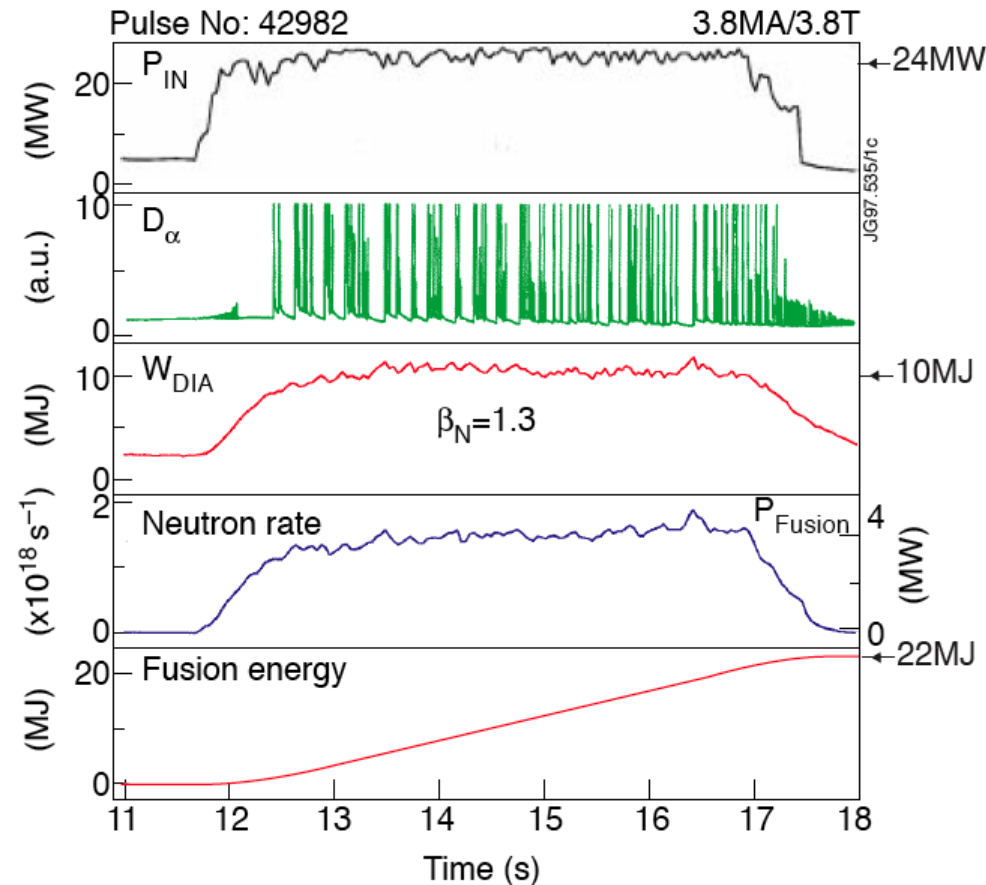




Fusion Power Production

JET ELMy H-mode

- Stationary fusion power production can be established in an ELMy H-mode - here 5s:
 - $P_{fus} = 4\text{MW}$
 - $Q = 0.16$
 - $E_{fus} = 22\text{MJ}$
- The duration of fusion power production in this type of plasma is limited to 5s by choice of the experimentalists:
 - duration of 5s allows stationary conditions to be demonstrated



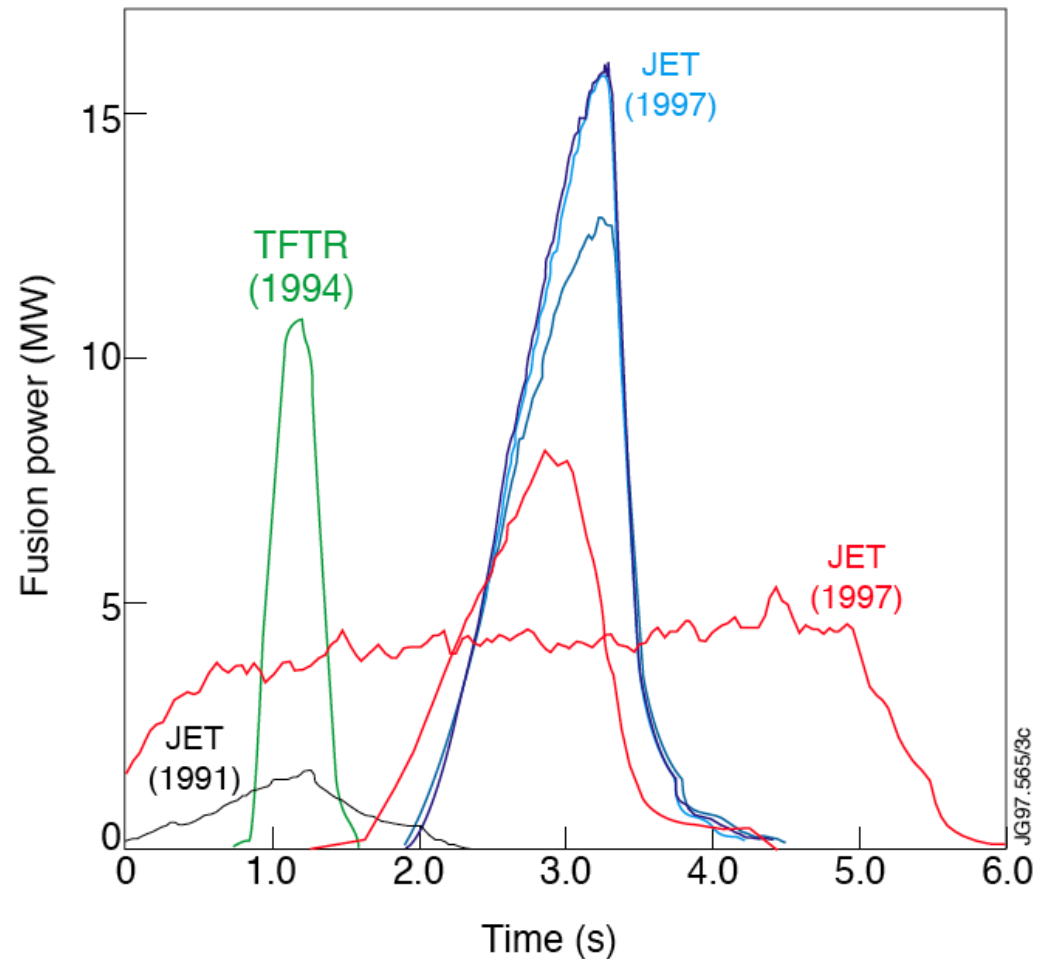
⇒ ITER reference plasma

J Jacquinot et al, Nucl Fusion **39** 235 (1999)



Fusion Power Production Overview

- Summary of best fusion power performance achieved in DT experiments in JET and TFTR





Conclusions

- **The magnetic fusion programme made the significant transition to operation with DT fuel mixtures during the 1990s**
- **Experiments in JET and TFTR made significant contributions to the progress of fusion energy research:**
 - demonstrated of safe handling of tritium
 - developed understanding of influence of tritium on plasma performance
 - Characterized behaviour of plasmas with significant α -particle content
 - confirmed expected behaviour of α -particles \Rightarrow **demonstrated α -particle heating of plasmas**
 - demonstrated fusion power production and verified code predictions of plasma fusion performance
 - addressed tritium retention and studied means for removal of tritium from vacuum vessel
 - subsequently, safe remote maintenance of tokamak



References

- **JET Team** Fusion energy production from a deuterium-tritium plasma in the JET tokamak, *Nuclear Fusion* **32**, 187 (1992).
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- **Fisher, R. K. et al.** Measurements of fast confined alphas on TFTR, *Physical Review Letters* **75**, 846 (1995).
- **Taylor, G. et al.** Fusion heating in a deuterium tritium tokamak plasma, *Physical Review Letters* **76**, 2722 (1996).
- **Hawryluk, R. J.** Results from deuterium-tritium tokamak confinement experiments, *Reviews of Modern Physics* **70**, 537 (1998).
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- **Keilhacker, M. et al.** High fusion performance from deuterium-tritium experiments in JET, *Nuclear Fusion* **39**, 209 (1999).
- **Jacquinet, J. et al.** Overview of ITER physics deuterium-tritium experiments in JET, *Nuclear Fusion* **39**, 235 (1999).
- **Thomas, P. R. et al.** Observation of Alpha Heating in JET DT Plasmas, *Physical Review Letters* **39**, 235 (1999).