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# Shape design of plasma facing components for stationary and transient power fluxes

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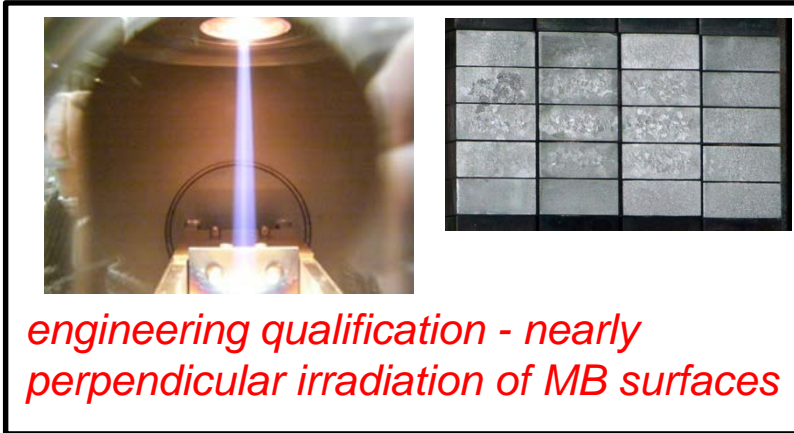
ITER International School

January 21-25, 2019

Daejeon, South Korea

*DISCLAIMER - The views and opinions expressed herein do not necessarily reflect those of the ITER Organization.*

## Full-W divertor from start of ITER operations



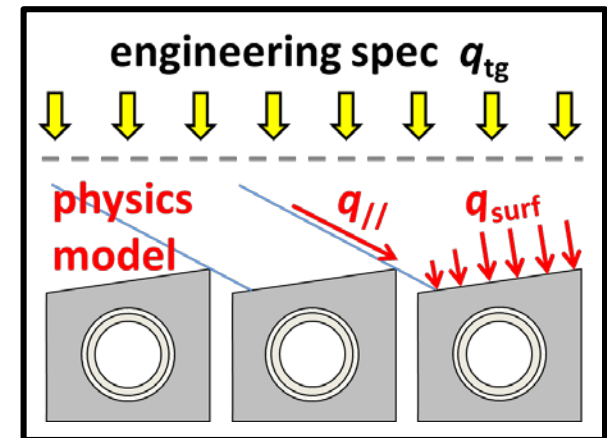
Heat load specifications prescribe maximum heat flux perpendicular to an ideal, axisymmetric divertor with no castellations or MB shaping.

The commonly heard phrase "steady state heat flux must be limited to 10 MW/m<sup>2</sup>" has its origin in such high heat flux tests.

-specific to ITER MB technology. Other technologies have different limits.

**Question: what will be the thermal response if we expose ITER MBs to a physics-based model of divertor plasma that delivers the specified power loads?**

**-near glancing B-field incidence angle ~3°; shaping; Larmor gyration around field lines**



J. P. Gunn, et al. , "Ion orbit modelling of ELM heat loads on ITER divertor vertical targets", Nuclear Materials and Energy (2017).

J. P. Gunn, et al., "Surface heat loads on the ITER divertor vertical targets", Nucl. Fusion 57, 046025 (2017).

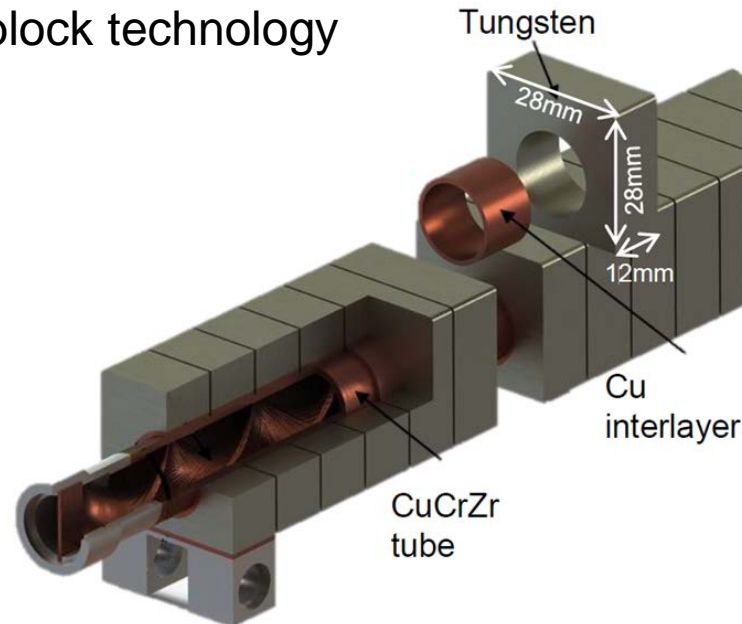
# HISTORICAL HEAT LOAD SPECIFICATIONS (FOR AN IDEAL AXISYMMETRIC DIVERTOR TARGET)

<b>Steady State (SS) inter-ELM detached regime</b>	<b>10 MW/m<sup>2</sup></b>	<b>to avoid W recrystallization</b>
<b>Slow Transient (ST) reattachment (300 events)</b>	<b>20 MW/m<sup>2</sup> → 10 s</b>	<b>to avoid critical heat flux (boil-out)</b>
<b>Fast Transient (FT) ELMs</b>	<b>~ 0.5 MJ/m<sup>2</sup></b>	<b>factor 2 margin against full surface melting of an initially cold monoblock</b>

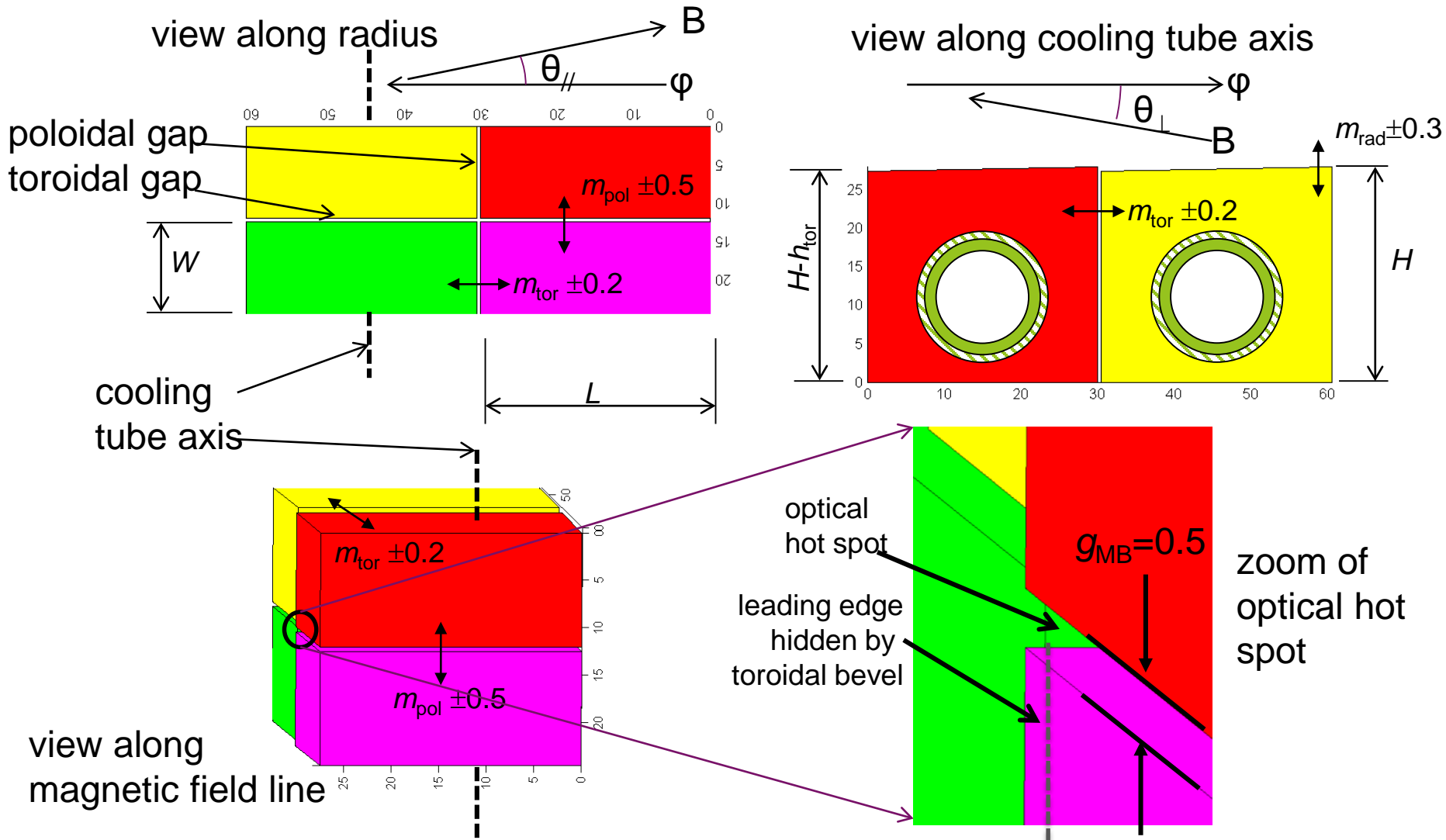
first part of talk

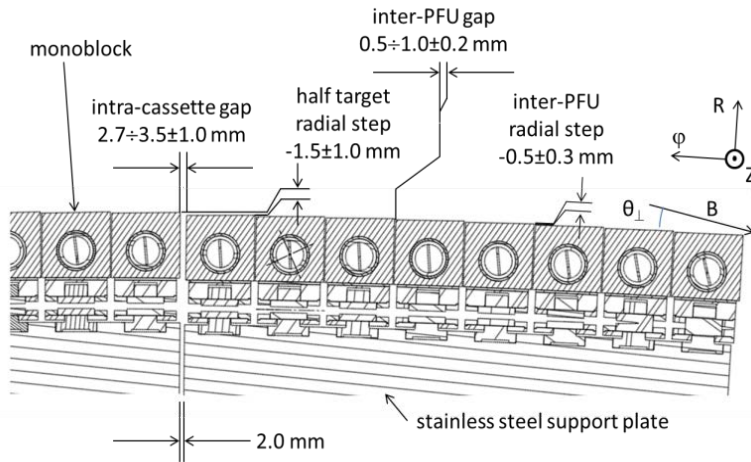
second part

ITER monoblock technology



# MONOBLOCK GEOMETRY AND B-FIELD ORIENTATION





Some of these tolerances have already been relaxed as a result of feedback from industrial suppliers, and they are complaining about others that are still too tight -consequences on divertor cost and performance

The studies reported here provide physics-based guidelines that give solid arguments for negotiations with suppliers

Literally thousands of 3D heat flux + thermal simulations were necessary to scan all tolerances and shaping alternatives

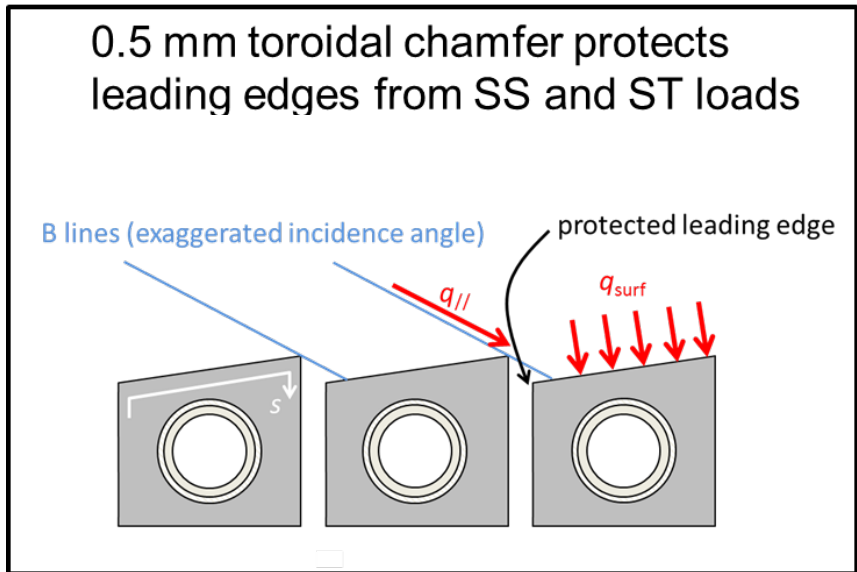
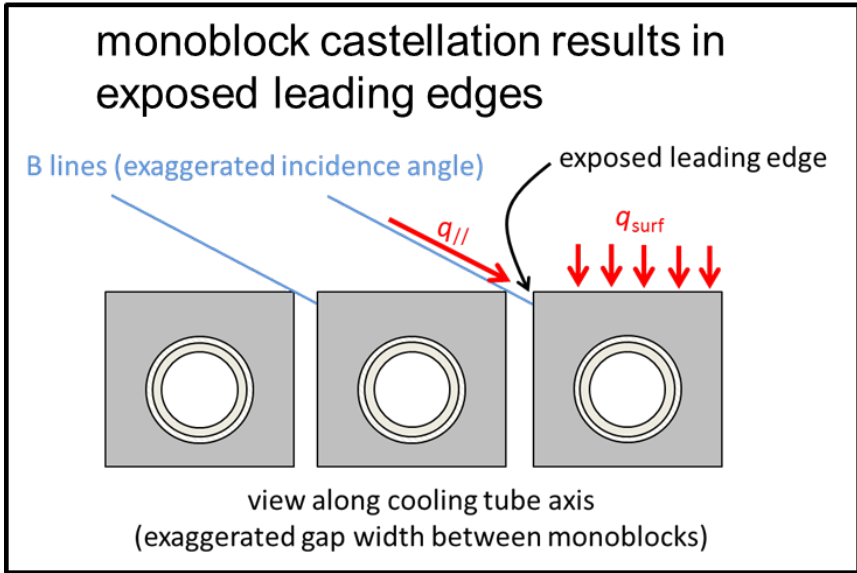
Message: good old analytic calculations and simple approximations remain a powerful tool - ANSYS is not God! Trust in your own brain.

feature	location	dimension [mm]	tolerance [mm]
gap	intra-PFU	$g_{MB} = 0.4$	IVT: $m_{pol} = +0.2$ -0.1 OVT: $m_{pol} = \pm 0.1$
	inter-PFU	IVT: $g_{PFU} = 0.5 \rightarrow 1.0$	$m_{tor} = \pm 0.2$
		OVT: $g_{PFU} = 0.5$	
	intra-cassette	IVT: $g_{PFU} = 2.7 \rightarrow 3.5$	$m_{tor} = \pm 1.0$
OVT: $g_{PFU} = 3.2$			
inter-cassette		$g_{PFU} = 20$	$m_{tor} = \pm 5$
radial step	intra-PFU	$\Delta r = 0.0$	$m_{rad} = \pm 0.3^*$
	inter-PFU	$\Delta r = -0.5$	$m_{rad} = \pm 0.3$
	intra-cassette	$\Delta r = -1.5$	$m_{rad} = \pm 1.0$
	inter-cassette	$\Delta r = -4.0$	$m_{rad} = \pm 2.0$
toroidal chamfer	both VTs	$h_{tor} = 0.5$	$\pm 0.1$

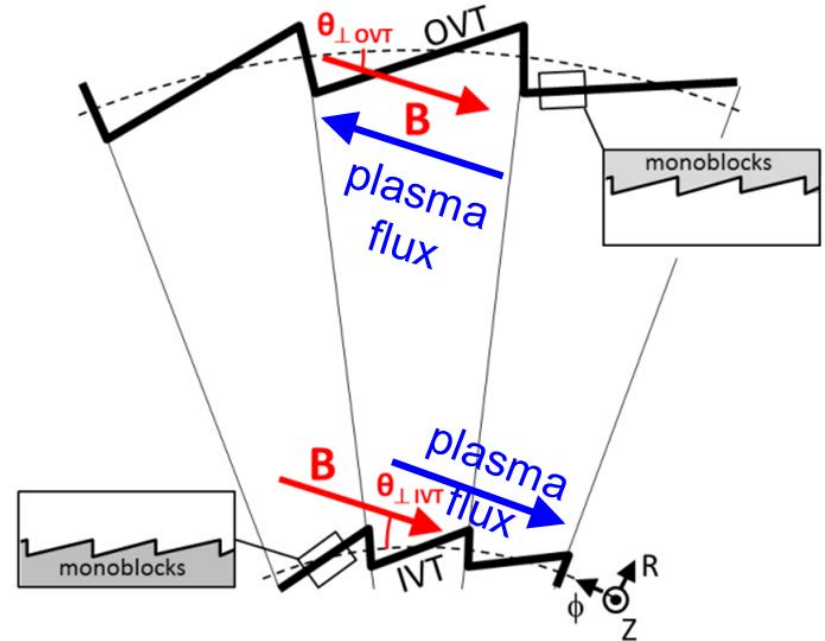
# Part 1

## inter-ELM

(i.e. "steady state")



schematic view of divertor illustrating target tilting and monoblock chamfer



target tilting and monoblock toroidal chamfer result in increased MB heat loads

heat flux delivered by plasma perpendicular to ideal target

$$q_{surf} \approx q_{tg} \frac{\theta_{\perp} + 0.5^{\circ} + 1^{\circ}}{\theta_{\perp}}$$

component tilting  
toroidal bevel

Percentage increase of plasma heat load

target	tilting+unshaped	tilting+bevel
IVT ( $\theta_{\perp}=3.2^{\circ}$ )	+16%	+47%
OVT( $\theta_{\perp}=2.7^{\circ}$ )	+19%	+56%

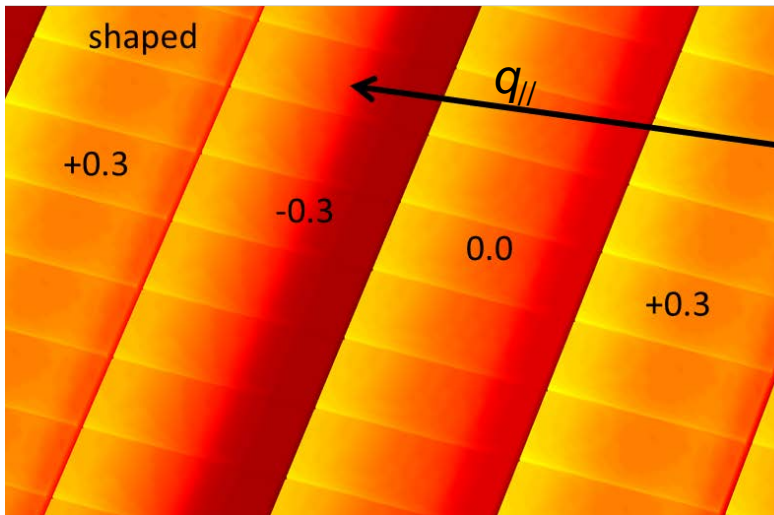
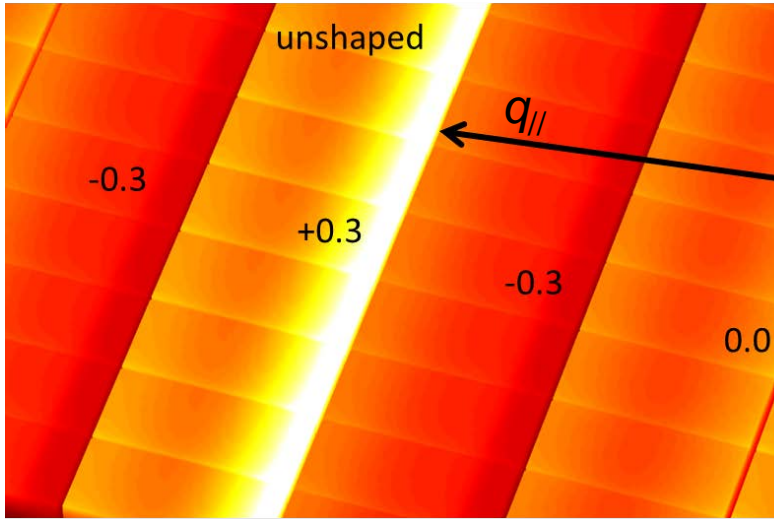
ST leading edge melting

No leading edge melting, but...  
SS recrystallization  
ST marginal surface melting  
FT ~90% surface melt threshold

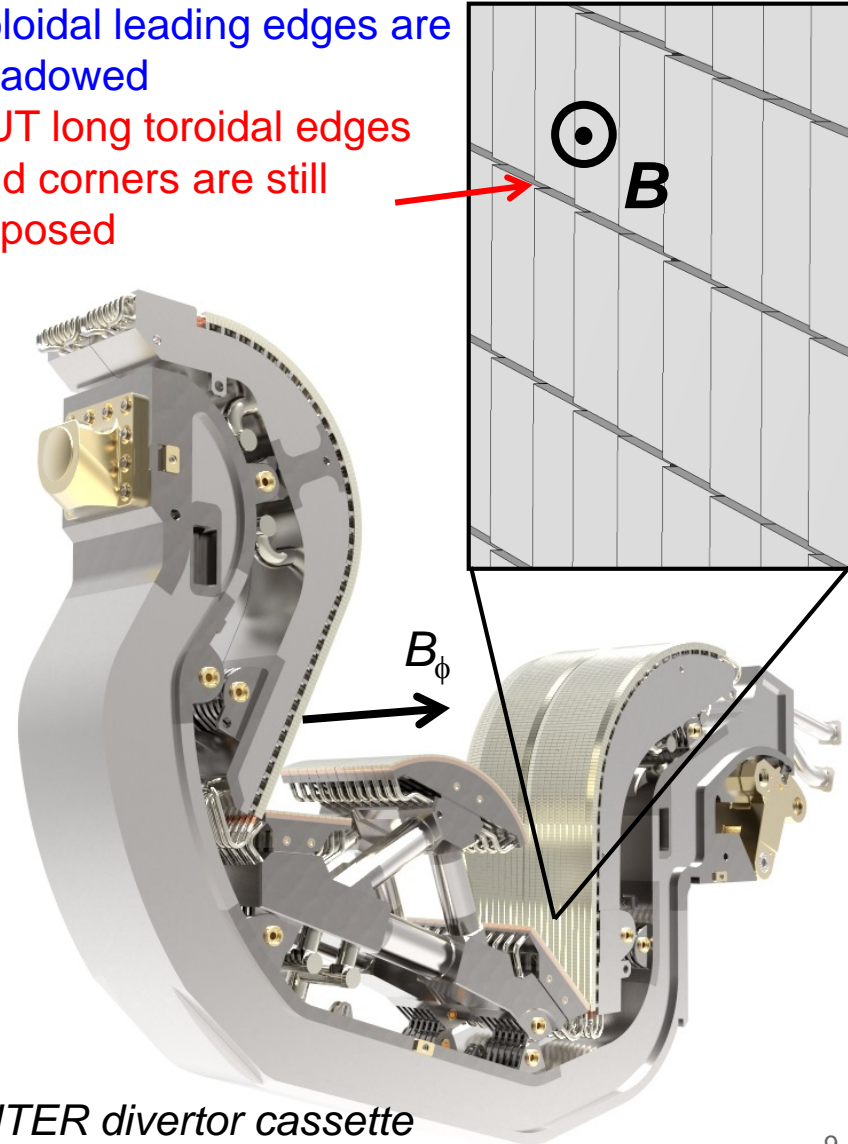


# POLOIDAL EDGES MOSTLY PROTECTED BY BEVELING: WHAT ABOUT TOROIDAL EDGES?

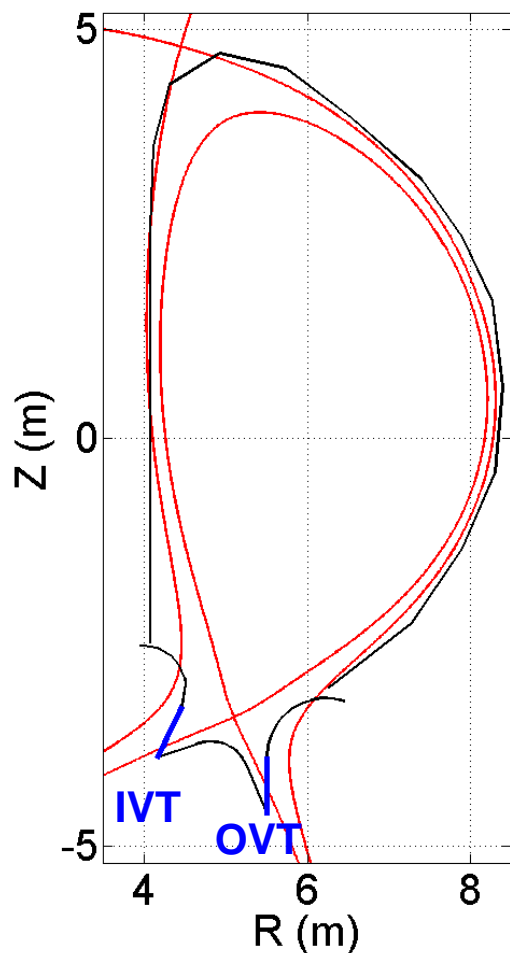
steady state thermal response of misaligned PFUs at outer vertical target



poloidal leading edges are shadowed  
BUT long toroidal edges and corners are still exposed



15 MA burning plasma



$P_{\text{SOL}} = 100 \text{ MW}$

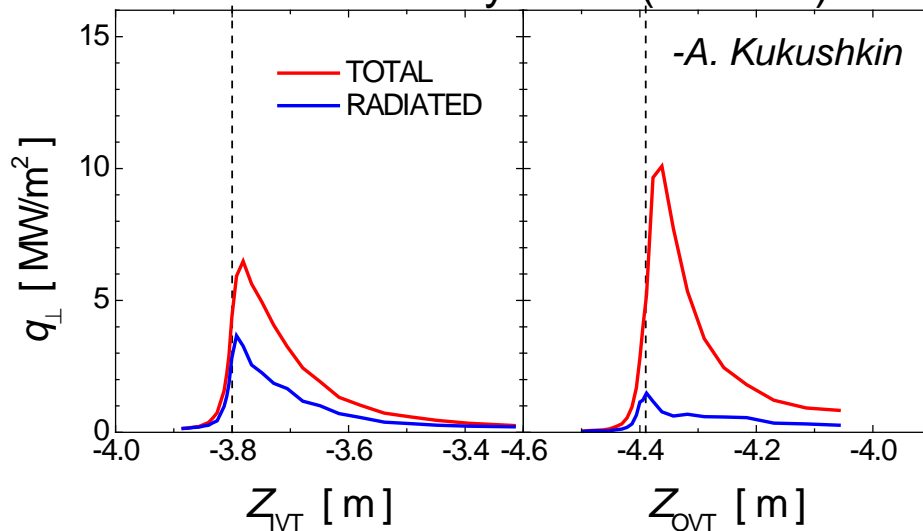
~2/3 to OVT

~1/3 to IVT

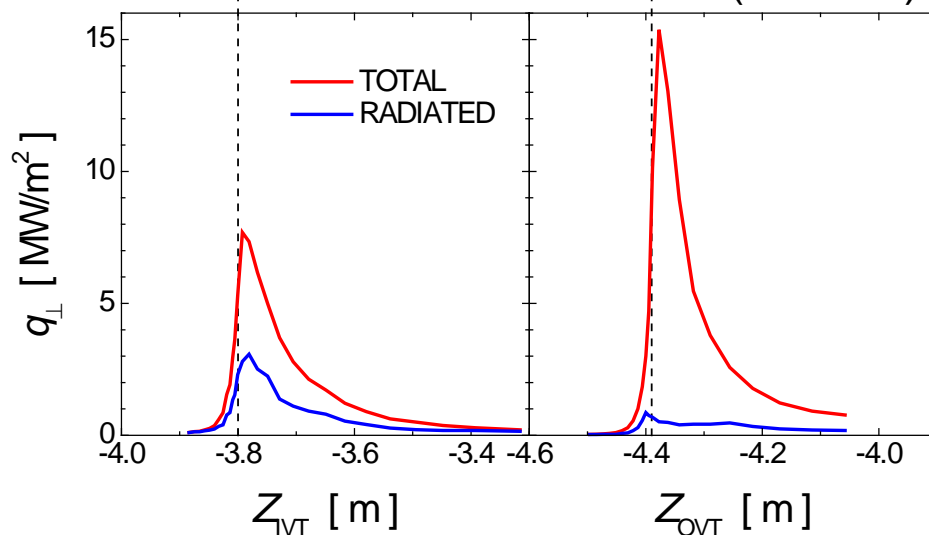
power  
dissipation by  
Neon injection

total power  
flux to divertor  
=  
plasma +  
photons +  
neutrals

nominal steady state (SOLPS)

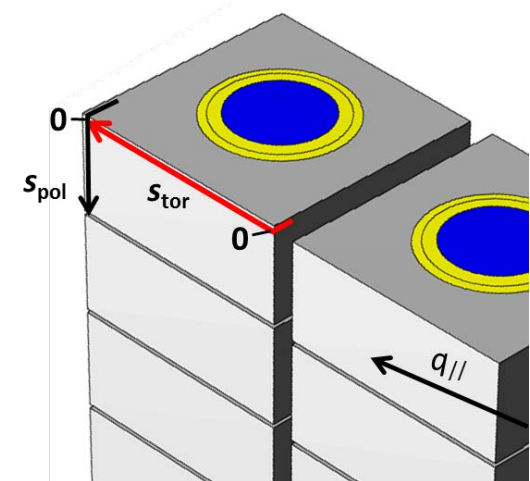
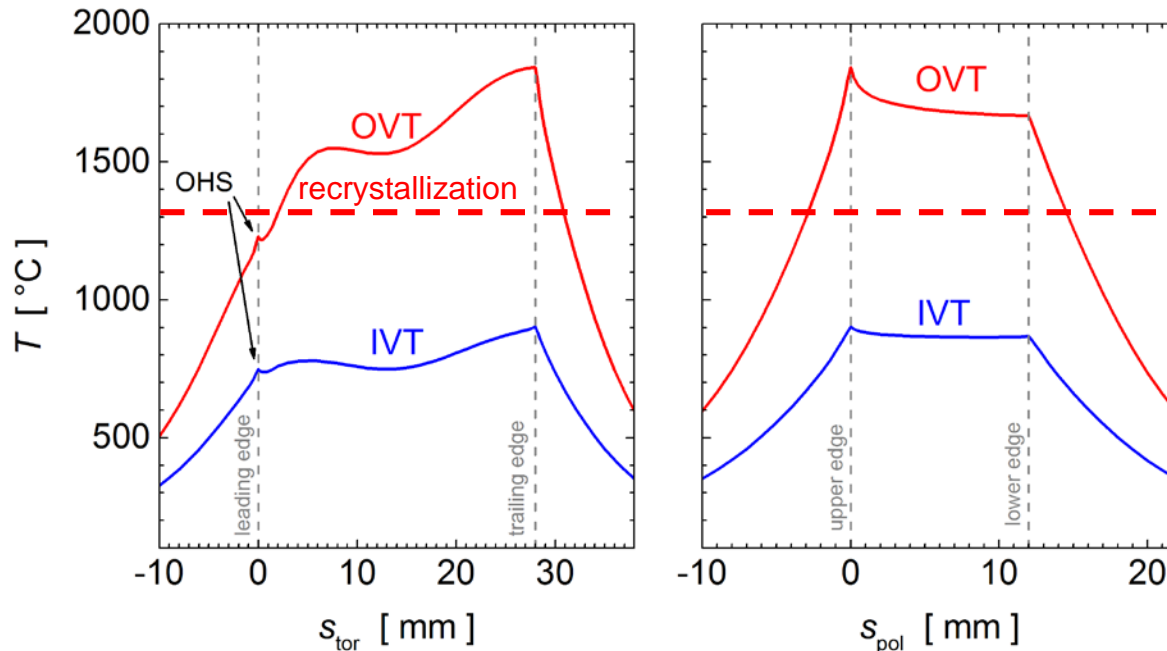
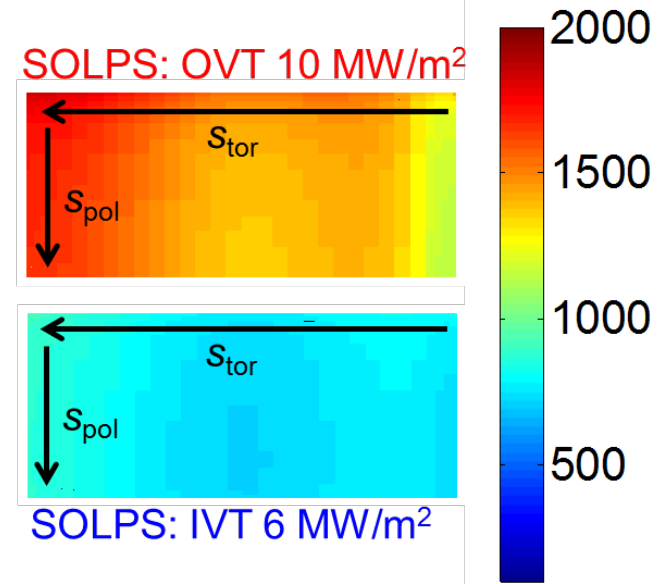


slow transient reattachment (SOLPS)



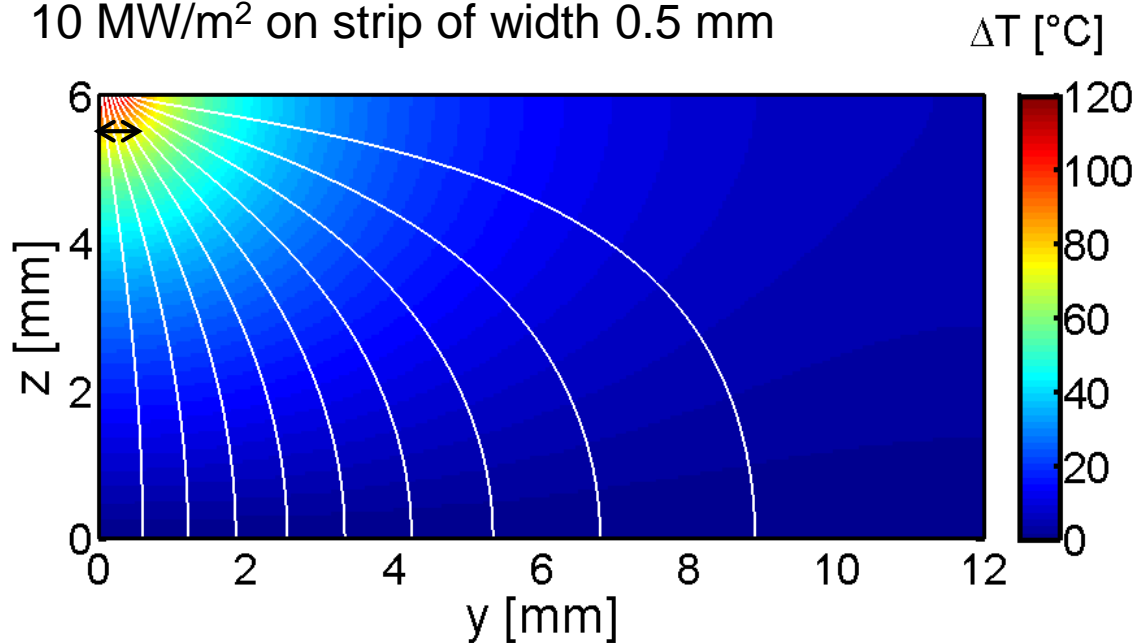
Long toroidal edges heat up  $\sim 100^\circ\text{C}$  more than top surface due to plasma entering toroidal gaps

target	$q_{tg}$ [ MW / m <sup>2</sup> ]	$q_{rad}$ [ MW / m <sup>2</sup> ]
IVT	6	3
OVT	10	1

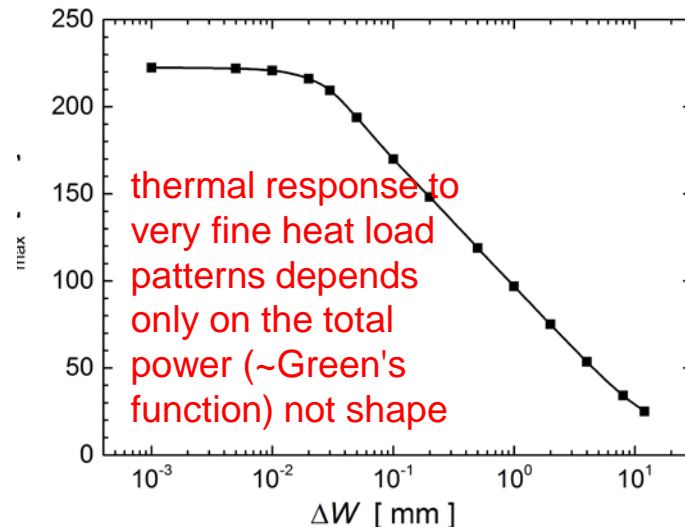


All the different heat sources can be decomposed and studied individually to understand the thermal response... Next slides

10 MW/m<sup>2</sup> on strip of width 0.5 mm



peak temperature vs width of strip (@ constant deposited power)



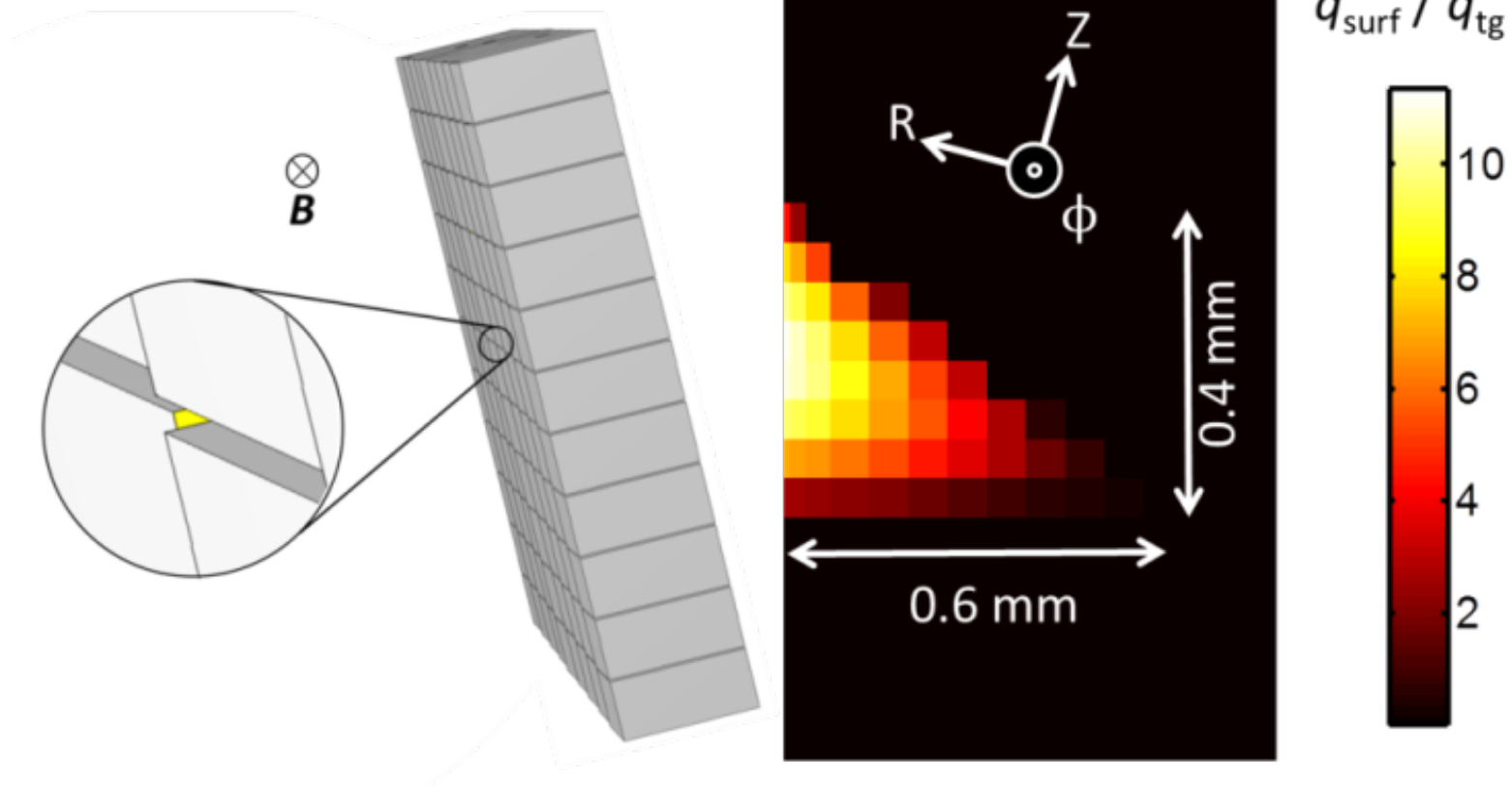
$$\Delta T(y, z) = \frac{-q_0}{\kappa} \left[ \frac{\Delta W}{W_{MB}} z + \sum_{n=1}^{\infty} \frac{2W_{MB}}{(n\pi)^2 \cosh \frac{n\pi H}{W_{MB}}} \sin \frac{n\pi \Delta W}{W_{MB}} \sinh \frac{n\pi z}{W_{MB}} \cos \frac{n\pi y}{W_{MB}} \right]$$

*solution by separation of variables*

Thermal properties of materials vary with temperature, but not dramatically, so linear approximation is valid (principle of superposition: the thermal response to multiple heat loads is the sum of the individual responses)

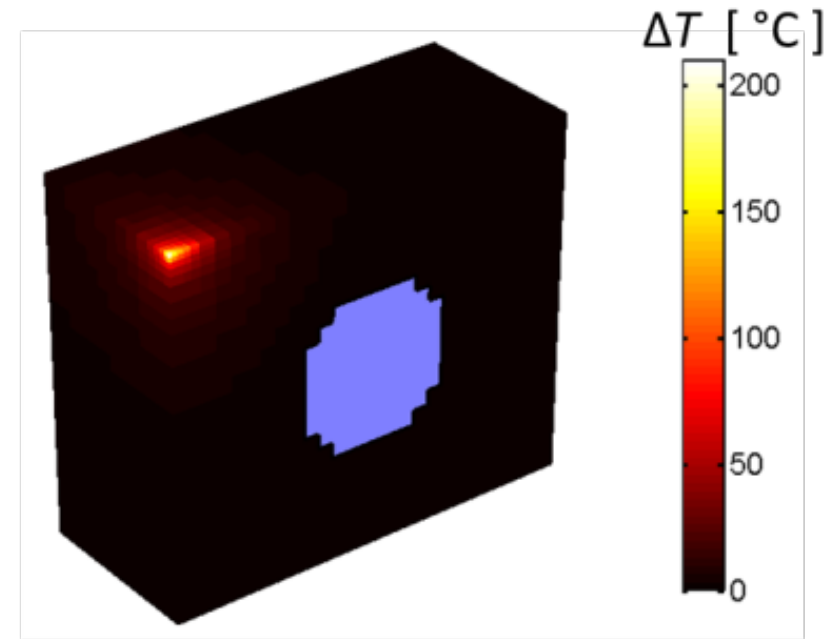
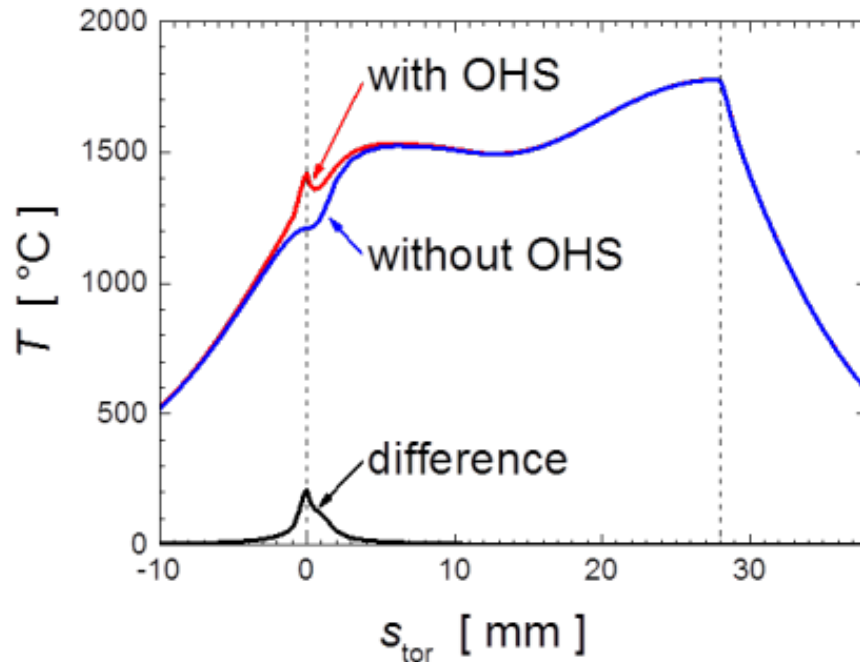
This is a 2D problem

- 1D linear source on boundary of 3D volume
- heat spreads in 2D, so small temperature gradient



poloidal leading edge visible through gap crossings  
 -direct irradiation by parallel heat flux ( $\sim 200 \text{ MW/m}^2$  in steady state)

IVT  $q_{tg}=10 \text{ MW/m}^2$   $q_{rad}=0$



This is a 3D problem

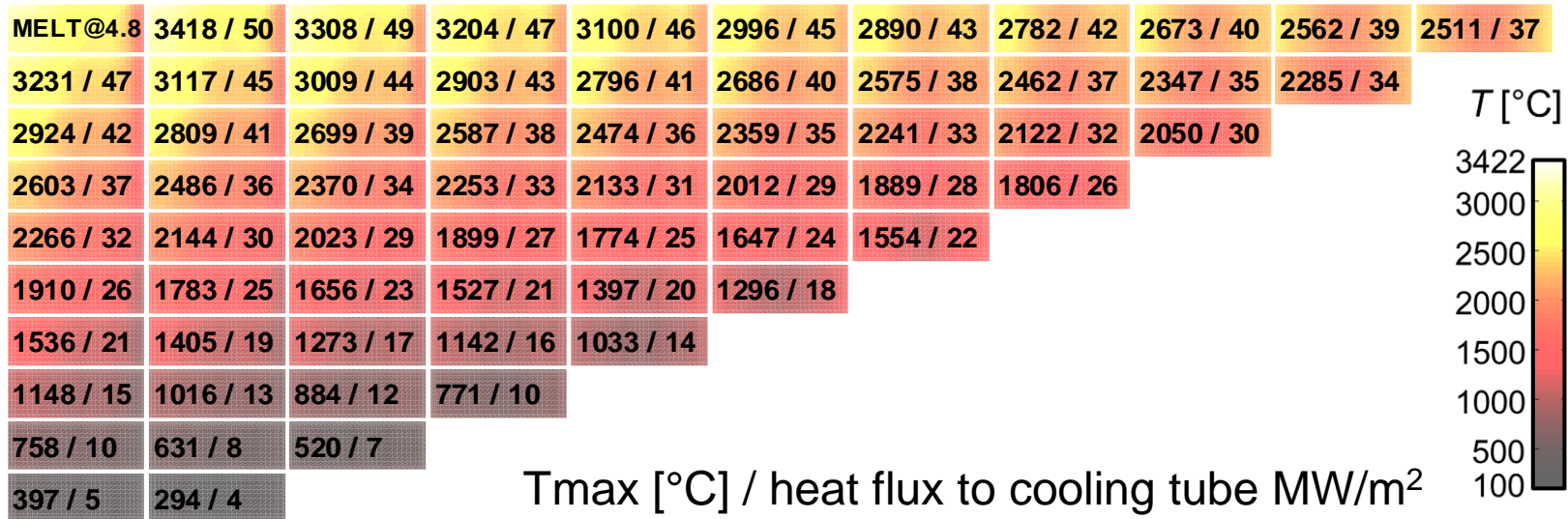
- 0D point source on boundary of 3D volume
- heat spreads in 3D, so small temperature gradient

N.B. temperature increase similar to hot strip, despite heat flux  $\sim 20X$  higher!  
We'll hear more about the OHS when we talk about ELMs later...

surface temperatures ~50% higher than high heat flux tests  
(because of tilt)

$q_{tg}$  [ MW/m<sup>2</sup> ]

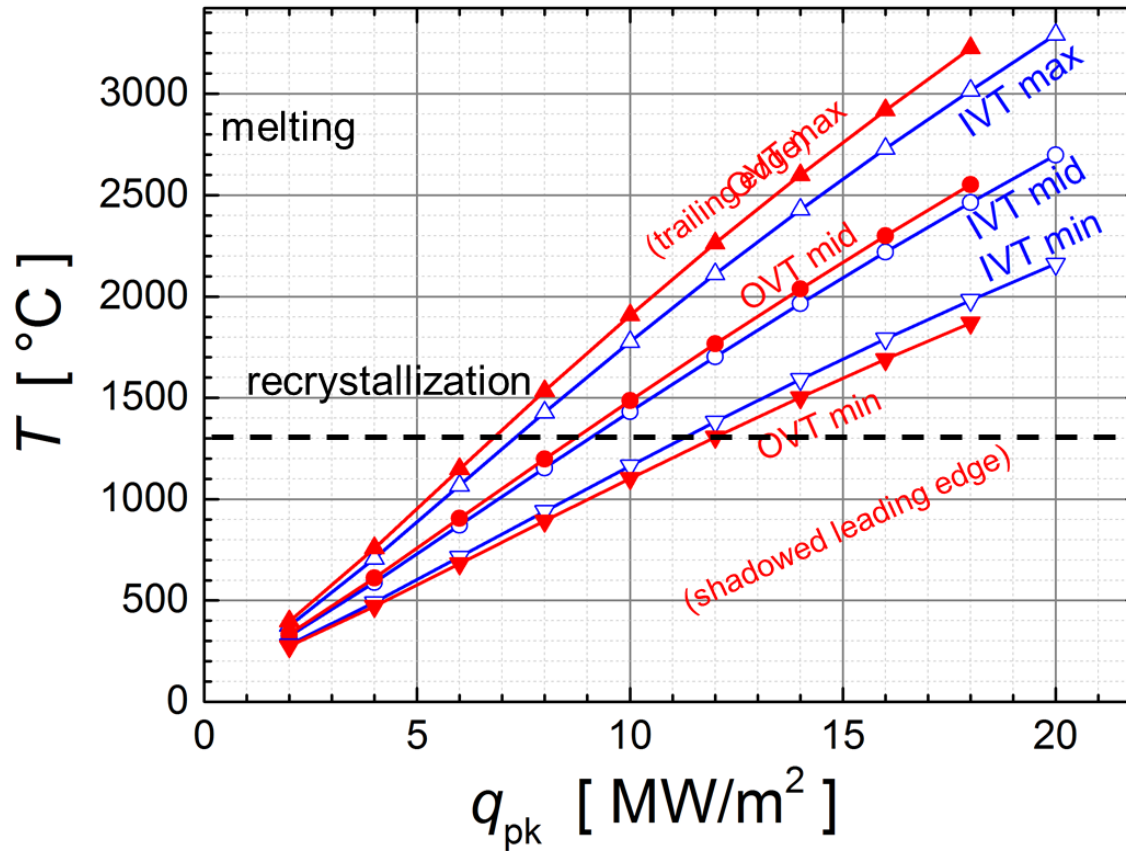
OVT intra-cassette



$q_{rad}$  [ MW/m<sup>2</sup> ]

\*assuming worst case misalignments

critical heat flux 40 MW/m<sup>2</sup> (formation of vapour layer, loss of heat handling, burnout)



Consequence of shaping - power flowing to divertor must be reduced  $\sim 2/3$  to avoid recrystallization

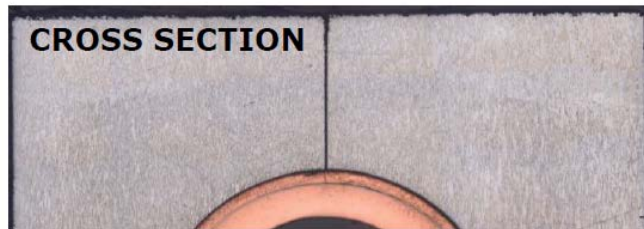


## inter-ELM loads

- shaping pushes surface temperature into recrystallization for steady state loads, and to marginal melting for slow transient loads (because of tilt)
- long toroidal edges heat up  $\sim 100^\circ\text{C}$  more than top surface (plasma flux into gaps)
- power to divertor would have to be reduced if recrystallization is to be avoided

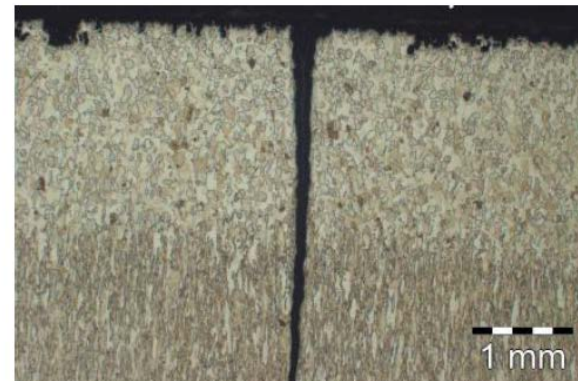
increase rate of Ne/N injection?

deeper detachment = loss of confinement (A. Huber, JET)



cracking of some W grades  
during slow transients

S. Panayotis (PSI Rome, 2016)

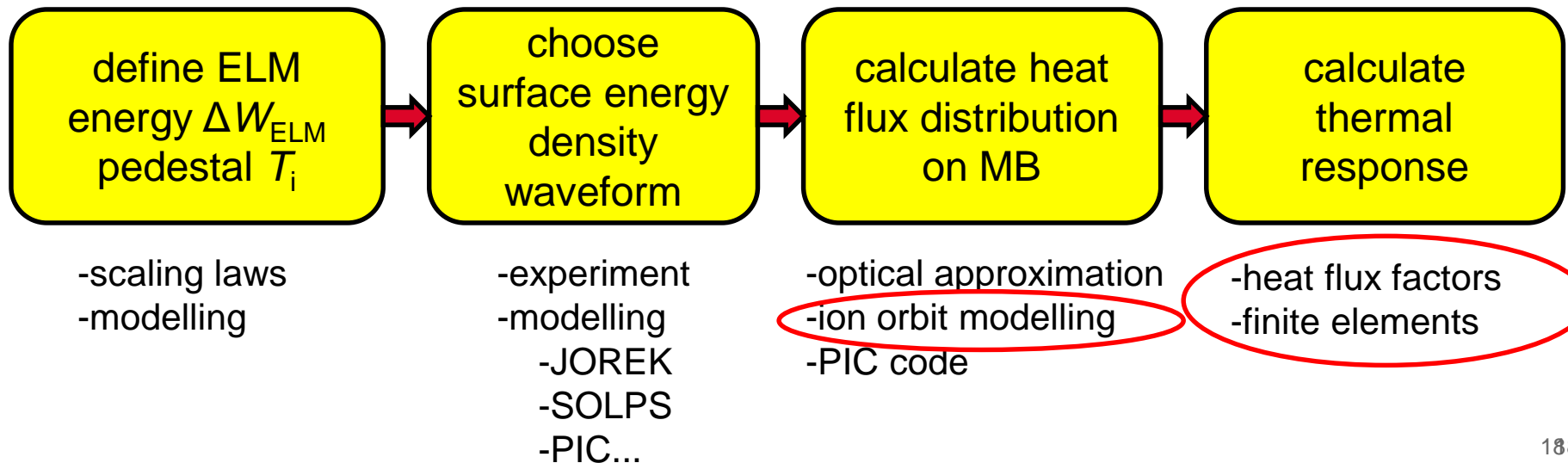


↑  
Recrystallized  
layer  
↓  
1 – 2 mm

G.Pintsuk, et al., SOFT2014

# Part 2

# ELMs



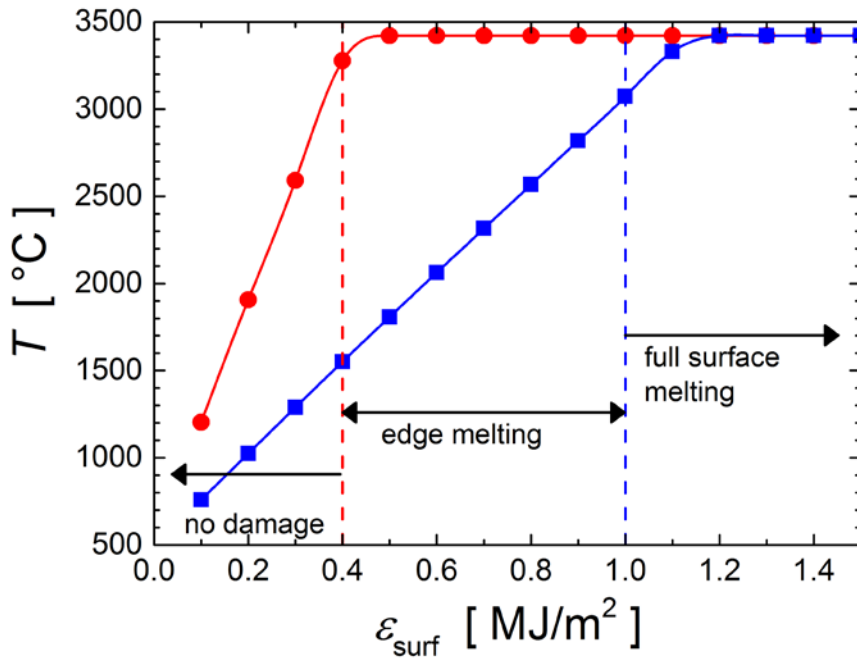
historical ITER limit  $\epsilon_{\text{surf}} \leq 0.5 \text{ MJ/m}^2$

-factor 2 margin against full surface melting (i.e.  $T_{\text{surf}} < 1700^\circ\text{C}$ )

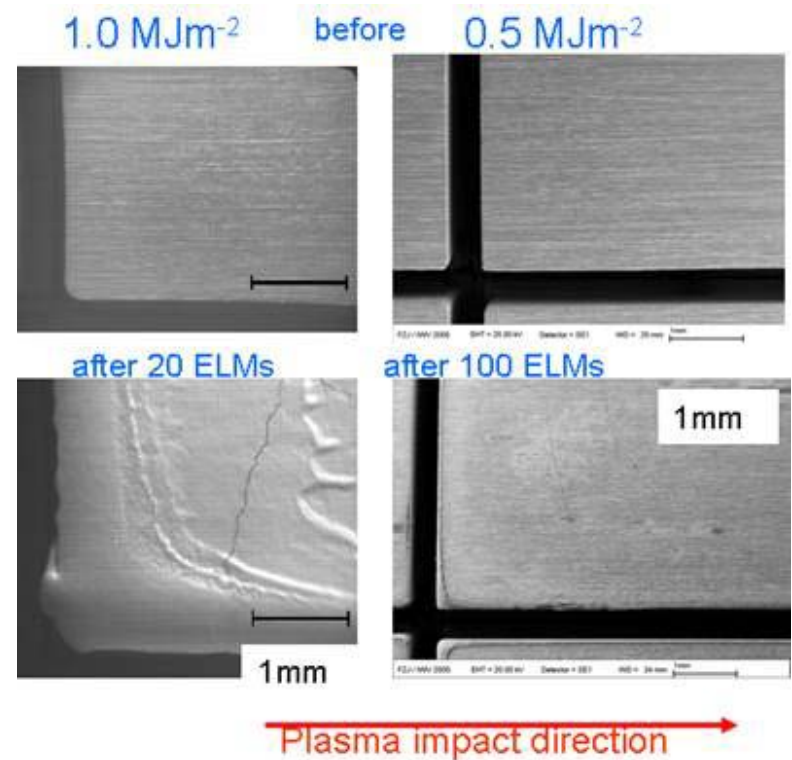
-marginal edge melting

N. Klimov, et al. JNM 390-391 (2009).

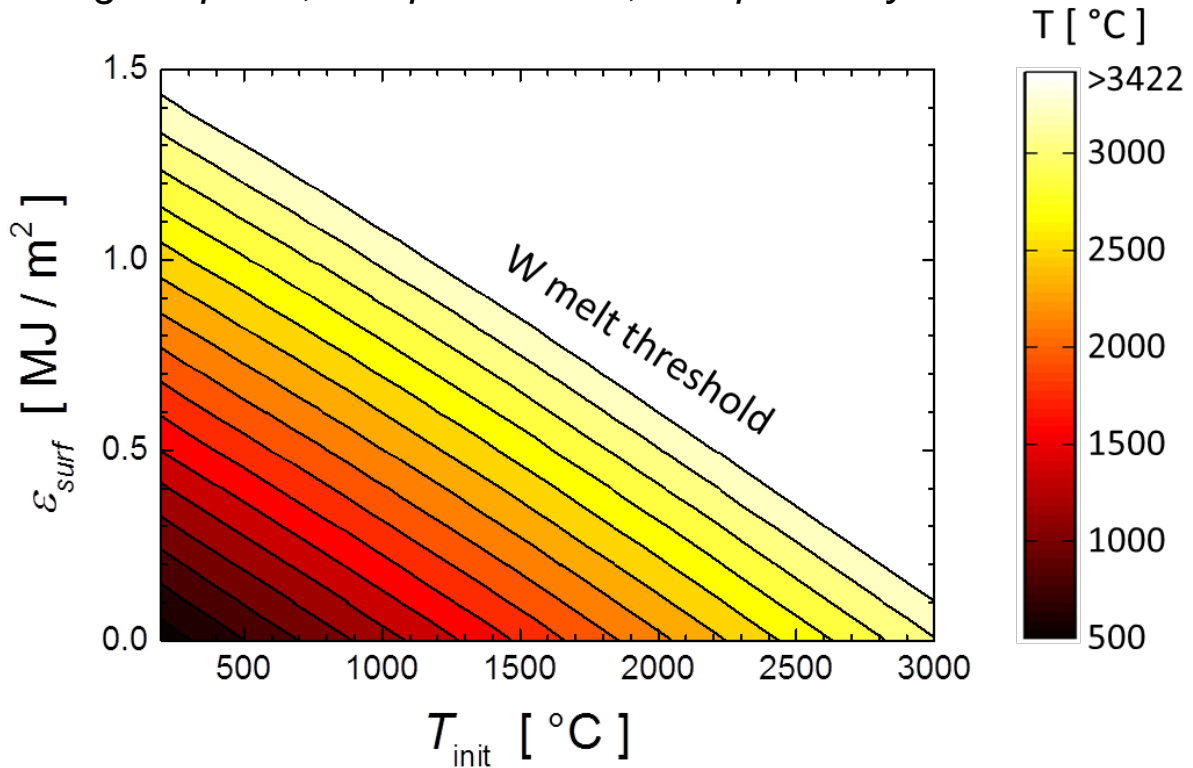
### QSPA square pulse



*data points from thermal model compared to visual evaluation of damage (dashed lines)*



*triangular pulse, 250  $\mu$ s rise time, 500  $\mu$ s decay time*



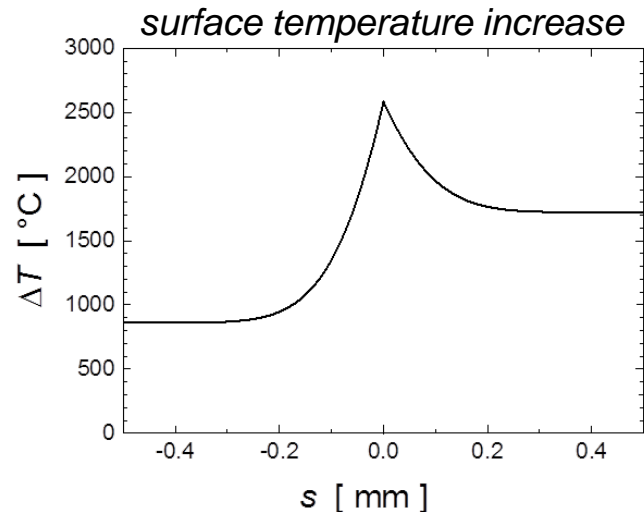
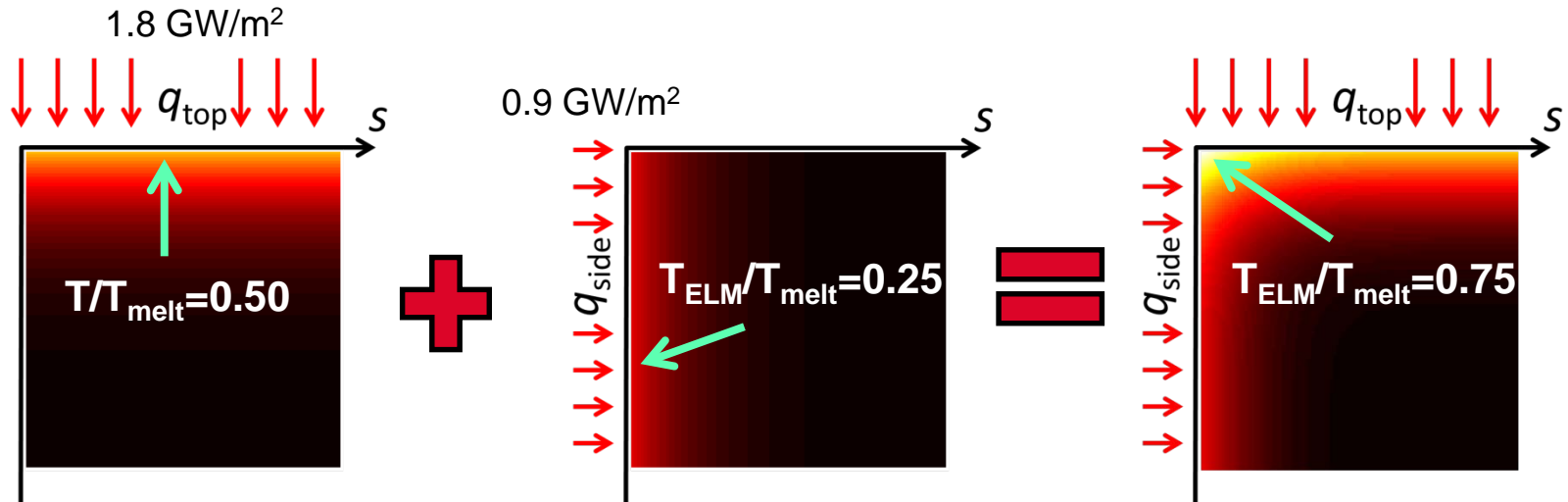
$\epsilon$  = ELM energy fluence  
i.e. total energy deposited  
during ELM event (time  
integral of heat flux)

historical ITER ELM limit ( $\epsilon_{tg}=0.5 \text{ MJ/m}^2$ ) generates temperature spikes  $\Delta T \sim 1100 \text{ }^\circ\text{C}$

**This factor 2 margin against melting is degraded for initially hot monoblocks**

*N.B. this limit applies to ideal, axisymmetric divertor with no castellations or shaping*

# AT A SHARP EDGE OR CORNER, THERMAL RESPONSE IS THE SUM OF 1D HEATING AT INDIVIDUAL FACETS



Exactly correct for linear case (temperature-independent thermal properties) (and 90° angles)  
Very good (<5%) approximation for non-linear (temperature-dependent thermal properties)

- 1) For a given magnetic field angle and specified ELM energy density, we calculate the corresponding  $q_{||} = q_{tg} / \sin \alpha$
- 2) We then launch that  $q_{||}$  at the monoblocks and calculate the local heat flux at all the surfaces of shaped monoblocks + worst case misalignments by 3D ion orbit simulations.

(weighting based on floating, ambipolar sheath)

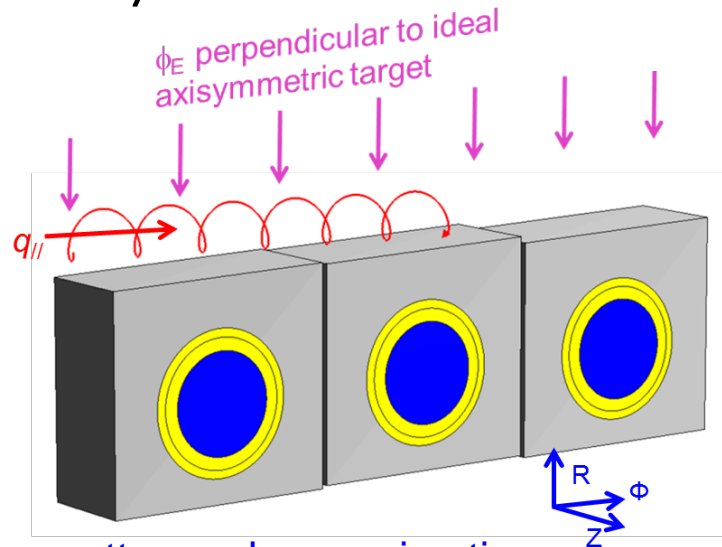
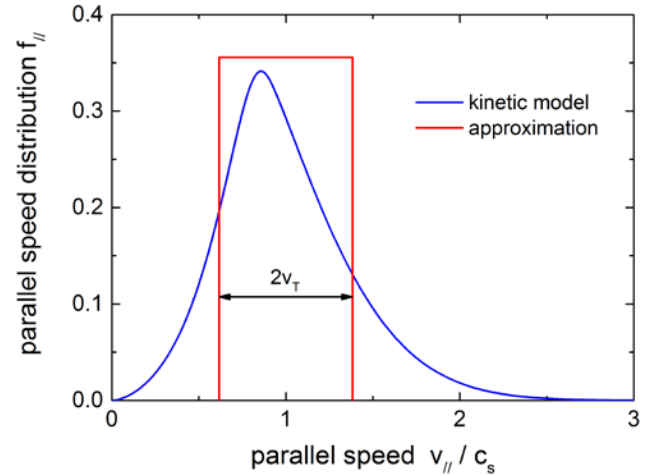
ion component:

- parallel speed distribution from kinetic model of SOL
- Maxwellian perpendicular speed distribution

electron component:

- optical approximation (field line tracing)

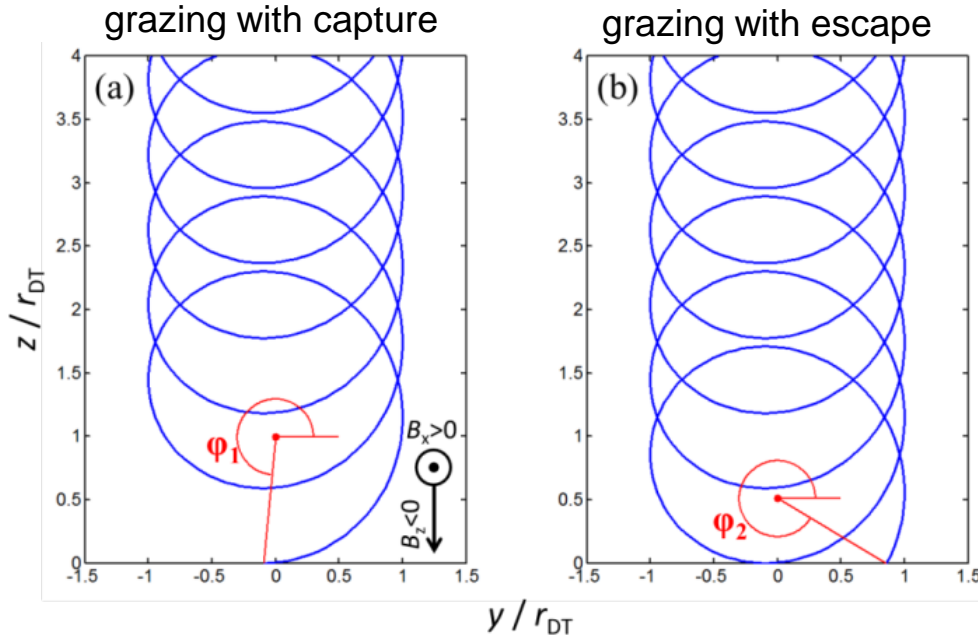
$$q_{surf} = \frac{5}{7} q_i + \frac{2}{7} q_e$$



Surprisingly (even to us) neglecting sheath E-fields is a pretty good approximation.

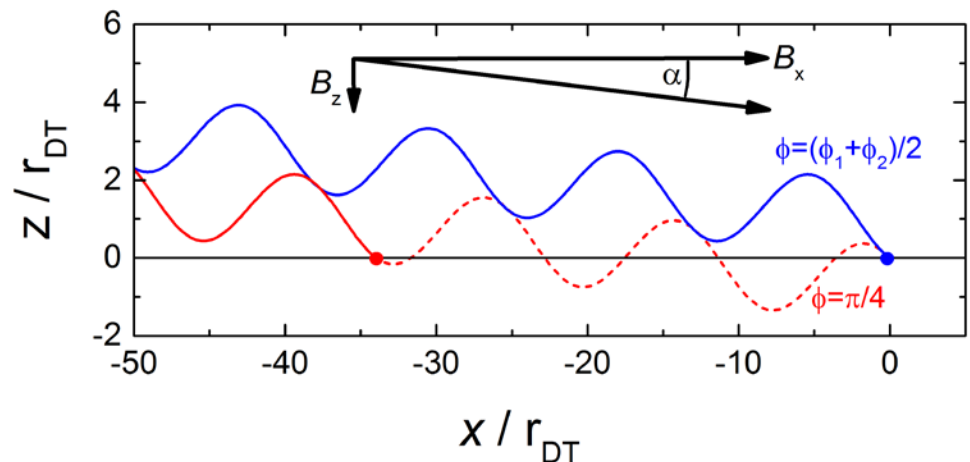
Confirmed by comparison with 2D particle-in-cell code SPICE

(M. Komm, et al., Nucl. Fusion **57**, 046025 (2017)).



Ions striking the surface have a restricted range of impact angles (nearly grazing)

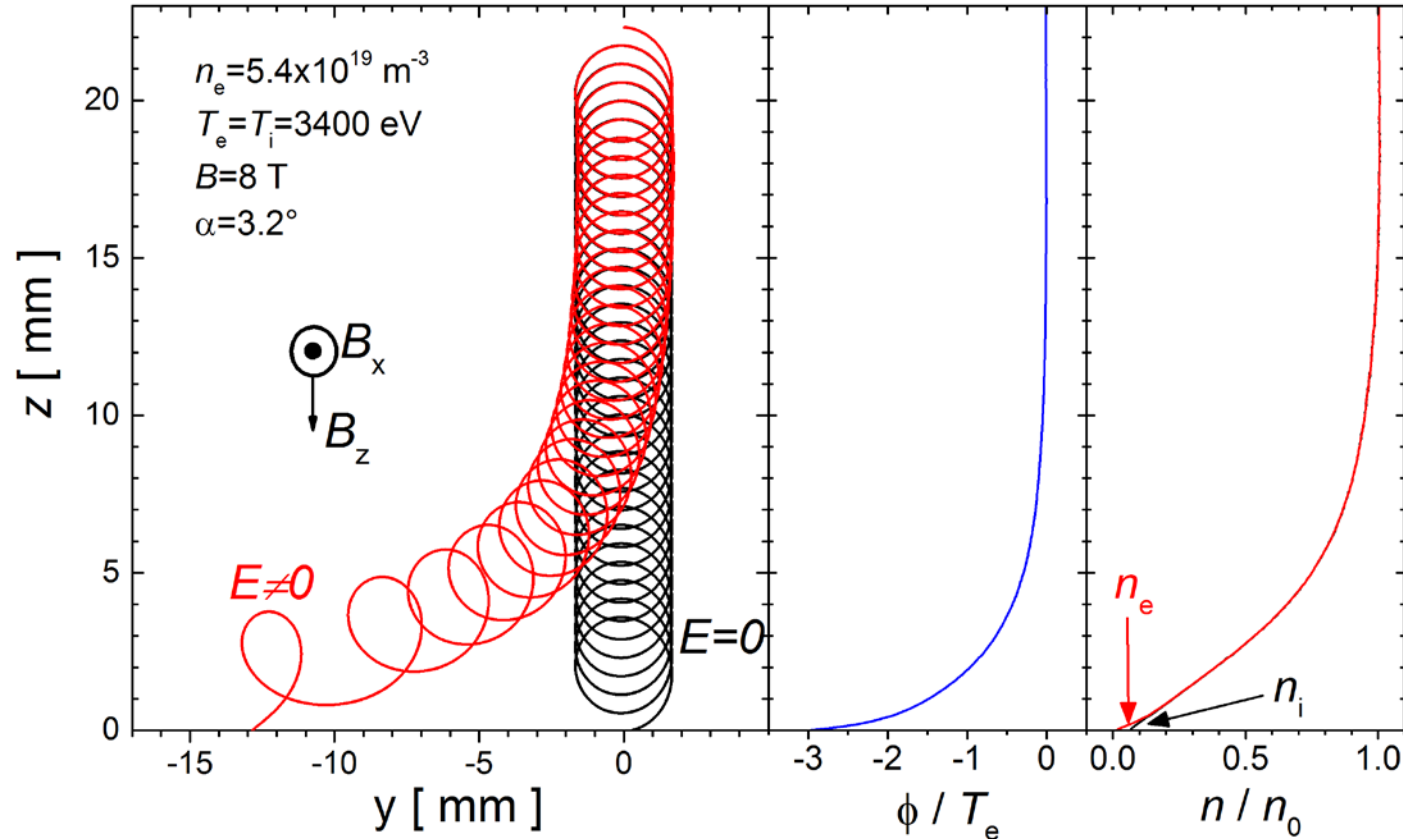
Angles outside this range do not exist because the ion would have struck the surface earlier



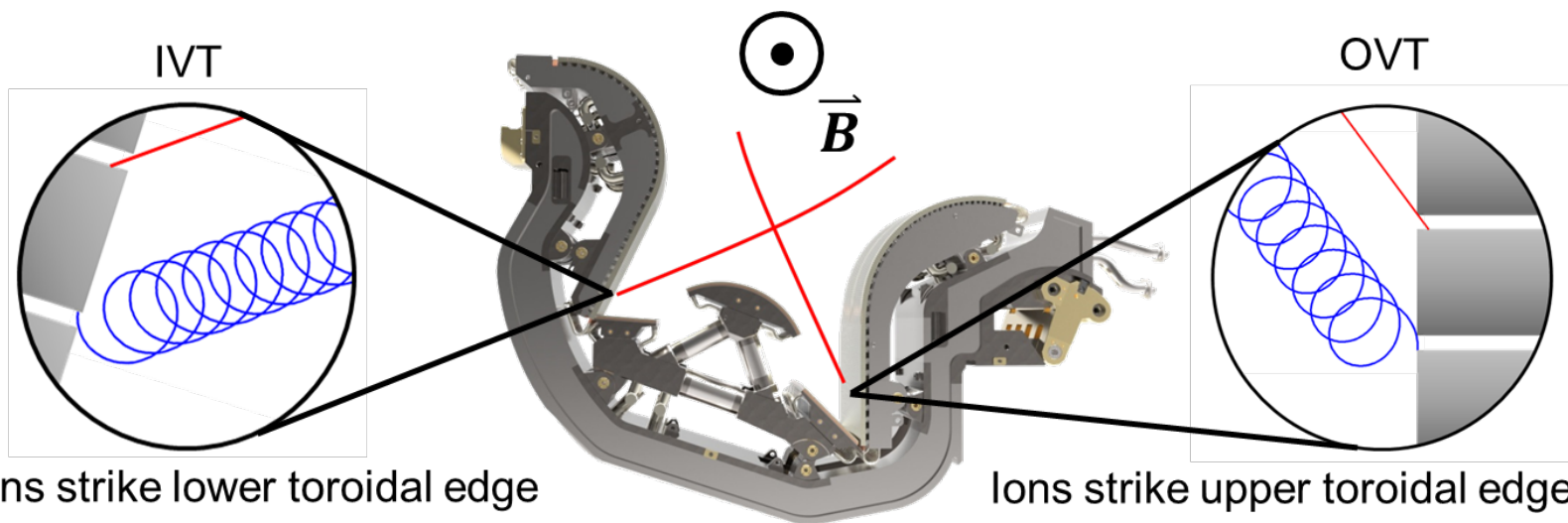
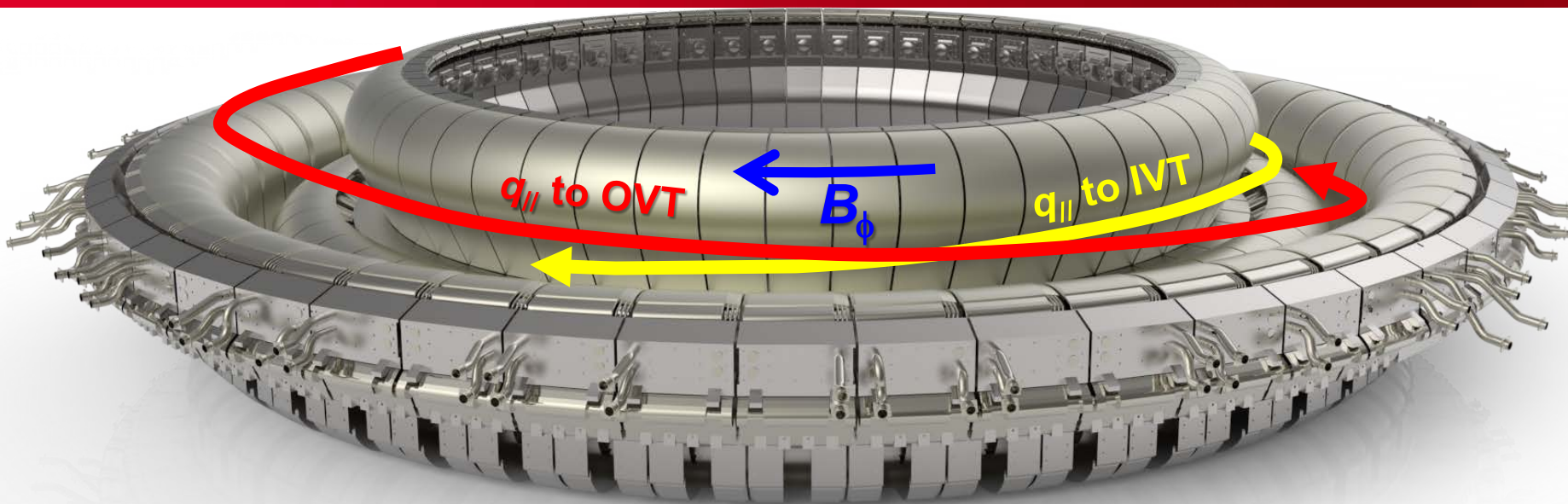
Electrostatic sheath (thin layer of strong electric field  $E \sim T_e / \lambda_D$ ) separates surfaces from plasma, keeping the plasma electrically neutral

Main effect is EXB drift parallel to surface - impact angles do not change much

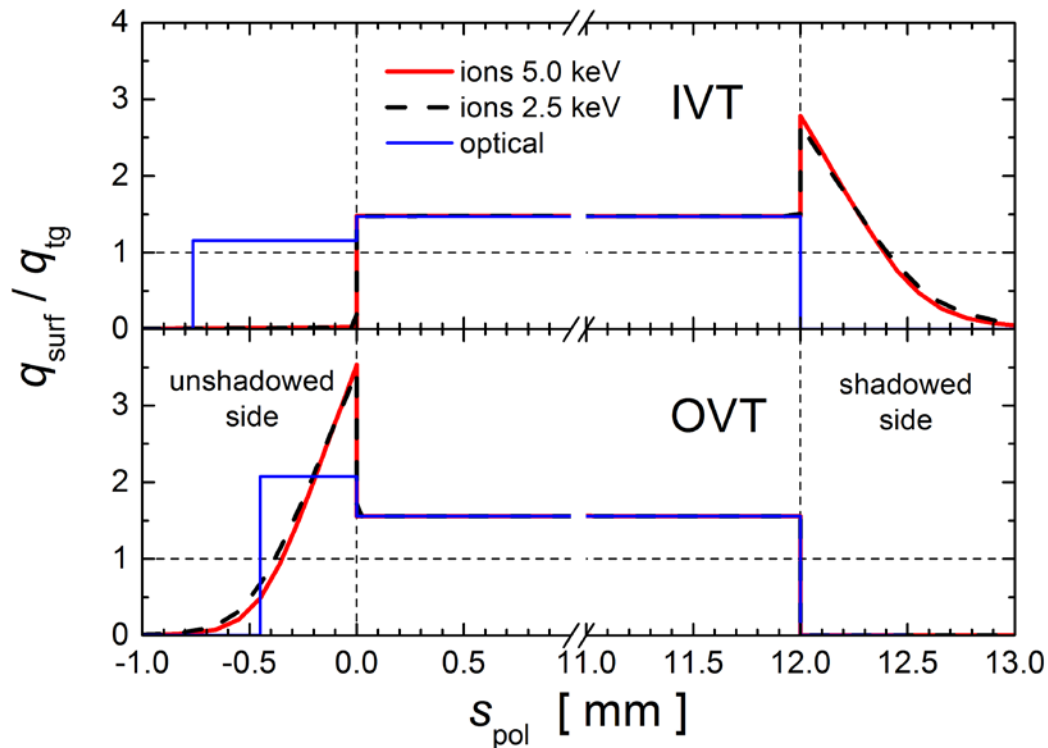
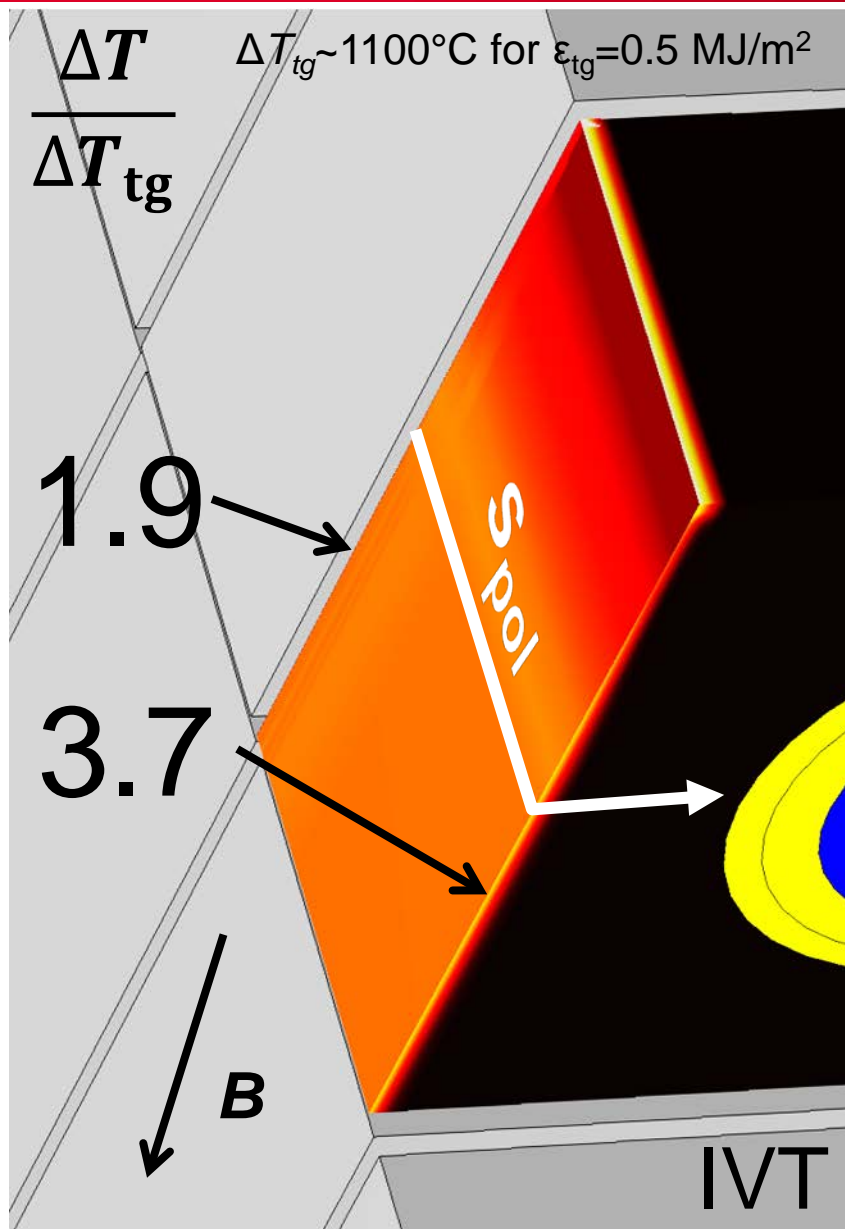
Assuming  $E=0$  seems dumb, but the approximation is "good enough"







electrons strike upper edges at both targets (tiny Larmor radius)



IVT: ions strike shadowed bottom side

OVT: ions strike wetted top side

-at both targets, electrons hit top side

First experimental confirmation of this asymmetry in COMPASS (for inter-ELM heat loads)

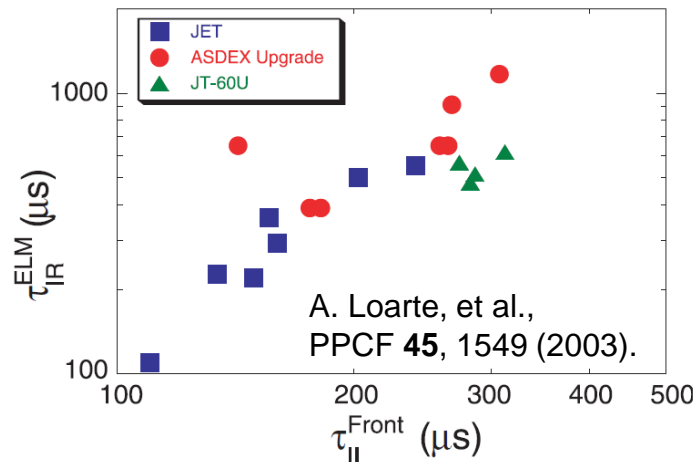
R. Dejarnac, et al., Nucl. Fusion **58**, 066003 (2018).

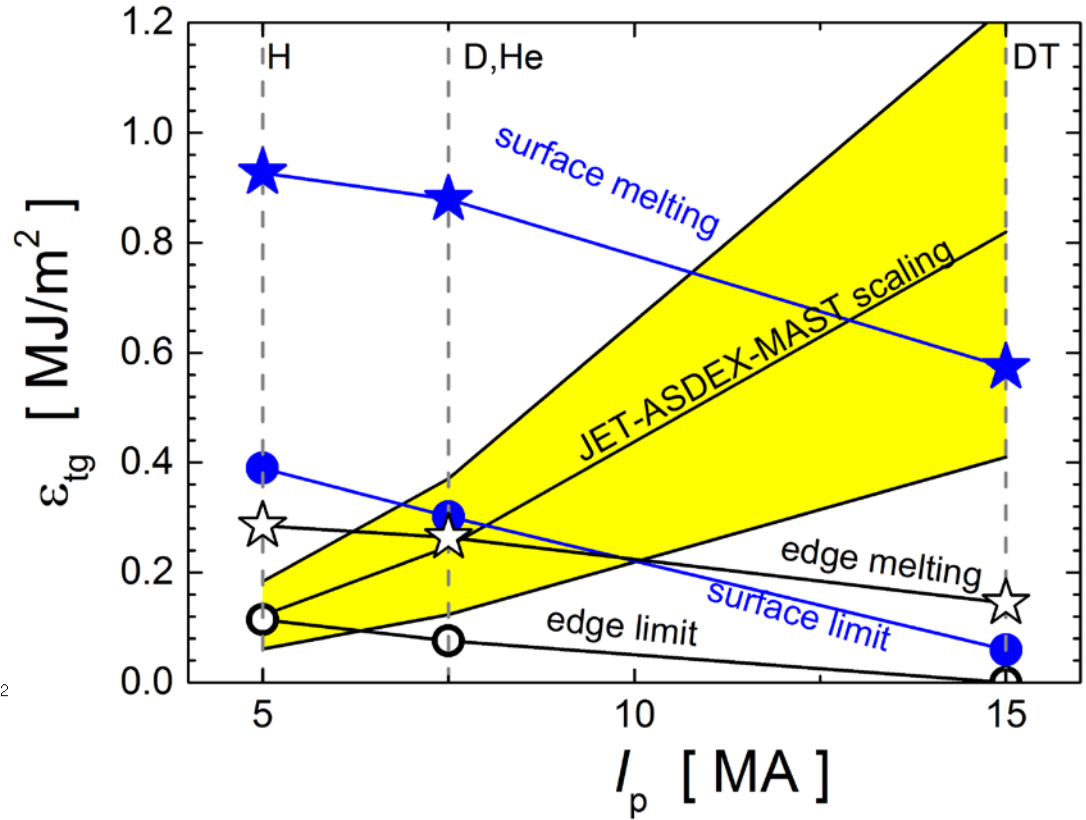
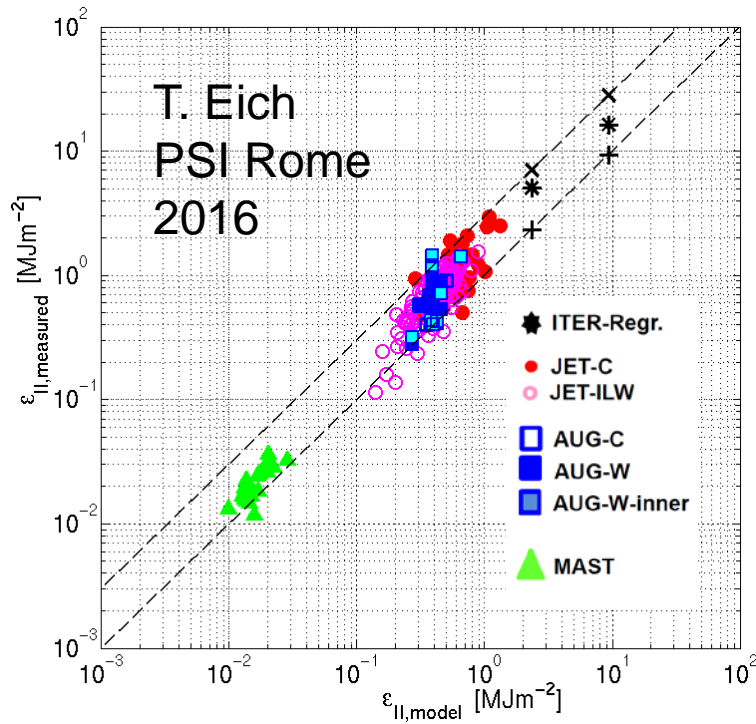
plasma	H	D or He	D+T
A/Z	1	2	2.5
$I_p$ [MA]	5.0	7.5	15
B [T]	1.76	2.65	5.3
$n_e$ [ $10^{20} \text{ m}^{-3}$ ]	0.3	0.4	0.8
$T_i$ [keV]	1.7	2.5	5.0
$\Delta t_{ELM}$ [ $\mu\text{s}$ ]	271	316	250
steady state $q_{tg}$ [ $\text{MW/m}^2$ ]	2.5	5	10
$T_{init}$ [ $^{\circ}\text{C}$ ] surface (edge)	450 (550)	800 (1000)	1500 (1900)

→ from inter-ELM thermal analysis with shaping

$$\Delta t_{ELM} = 250 \sqrt{\frac{2A}{ZT_i}} \quad [\mu\text{s}]$$

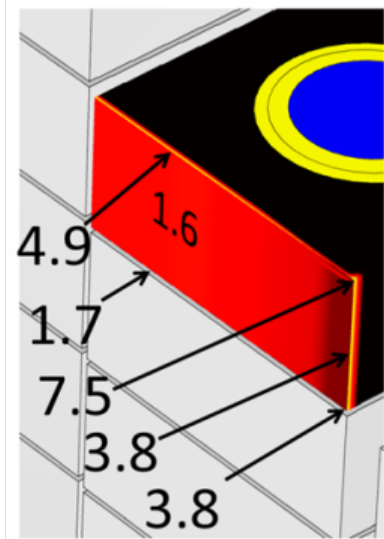
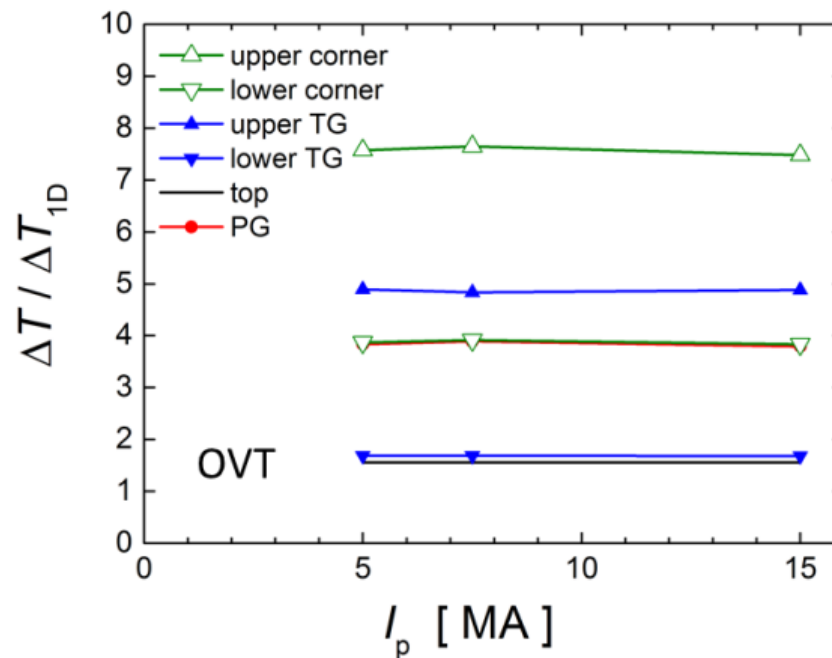
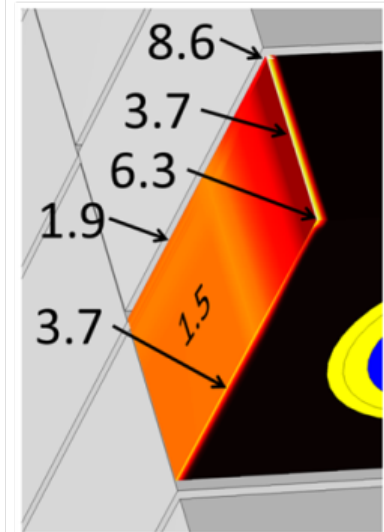
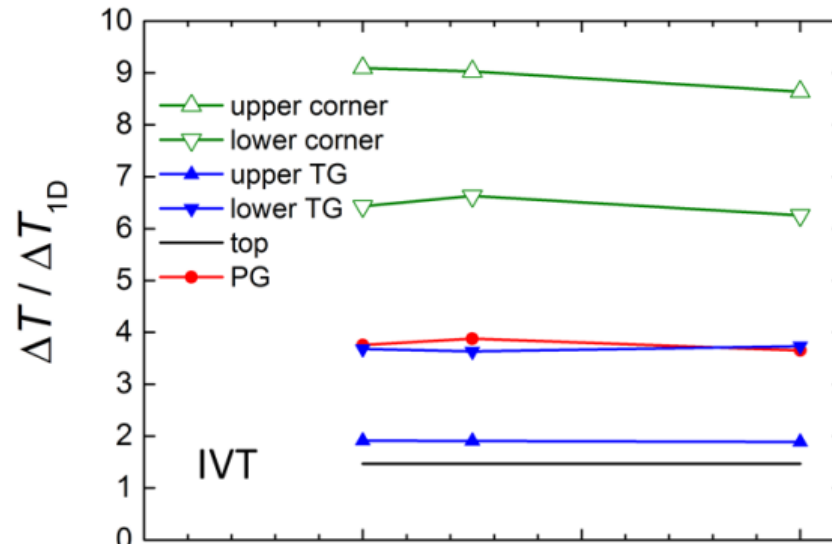
*ELM rise time:*  
*empirical scaling assuming free streaming from midplane to target at ion sound speed*





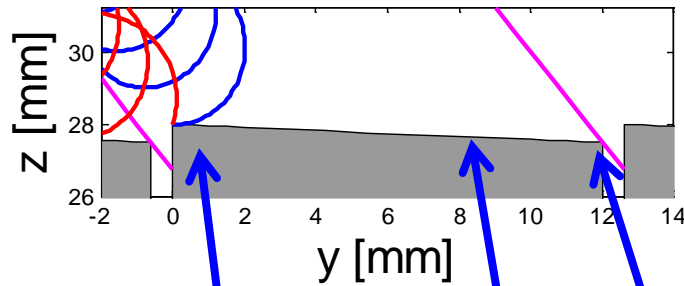
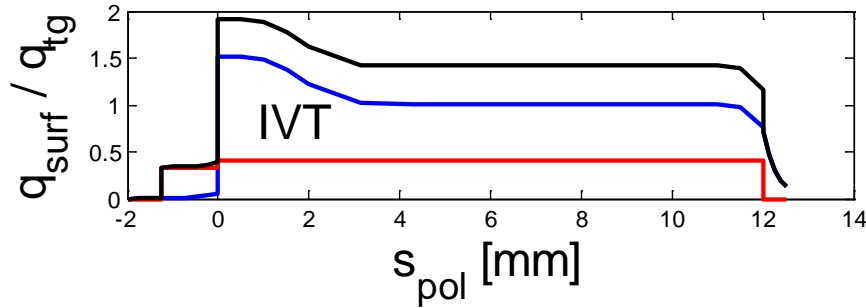
scenario	full surface melting?	edge melting?
pre-nuclear hydrogen 5MA	avoided with wide margin	avoided with narrow margin (less than 2)
pre-nuclear D or He 7.5 MA	avoided with narrow margin (less than 2)	possible during largest ELMs
DT nuclear burn 15 MA	unavoidable	unavoidable

Bonus:  
optical hot spot!  
heat load is  
sufficient to trigger  
tungsten *BOILING*  
at every ELM



At IVT, ions and electrons flow to opposite sides of the toroidal gap

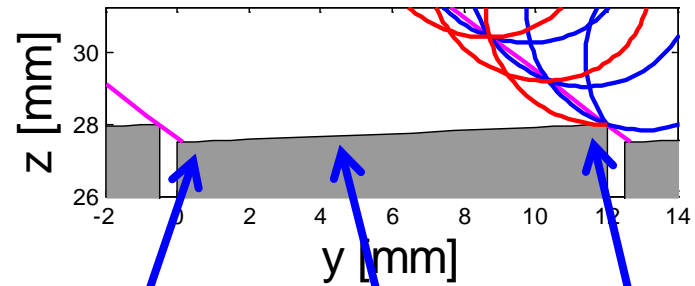
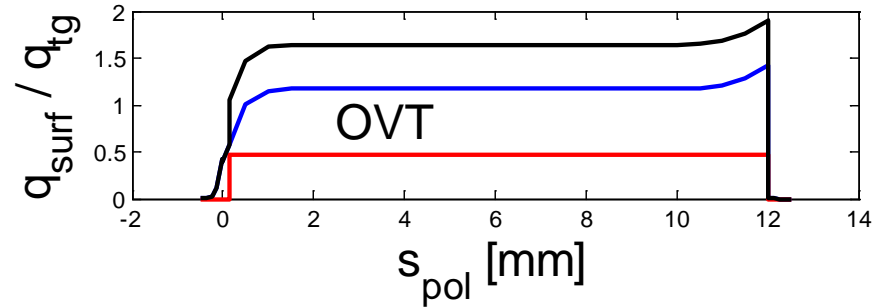
-poloidal beveling to protect against ELMs cannot fully succeed because either ions or electrons are affected, but never both



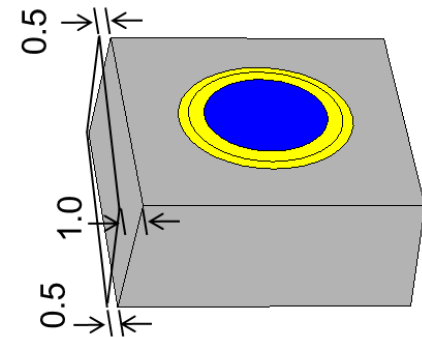
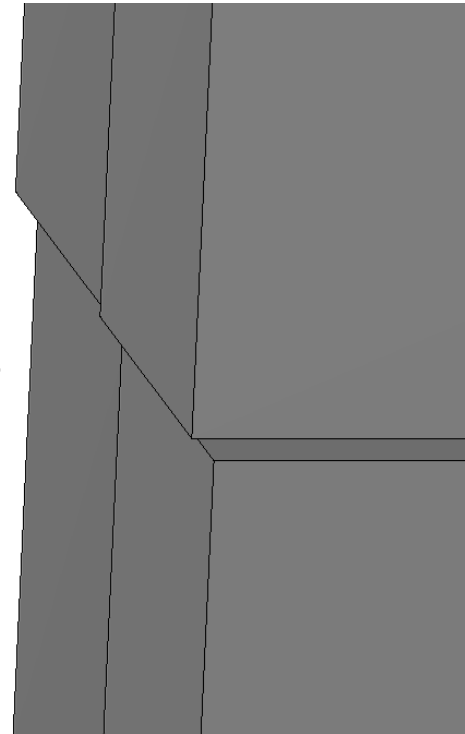
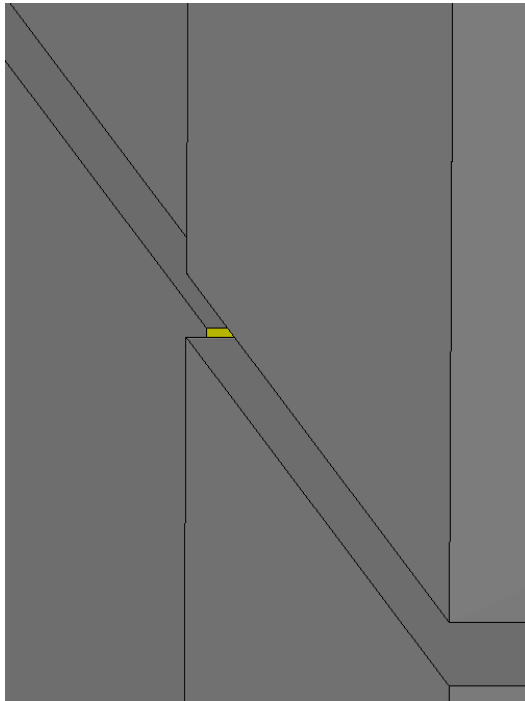
increased heating at upper edge  
decreased heating on main surface  
shadowing of lower edge from ELM ions

At OVT, both electrons and ions flow to the same side

Combined poloidal and toroidal bevels has the potential to mitigate the ELM and inter-ELM TG loading problem



shadowing of upper edge from both ion and electron loads  
slightly increased heating on main surface  
slight increase of ion heating at lower edge

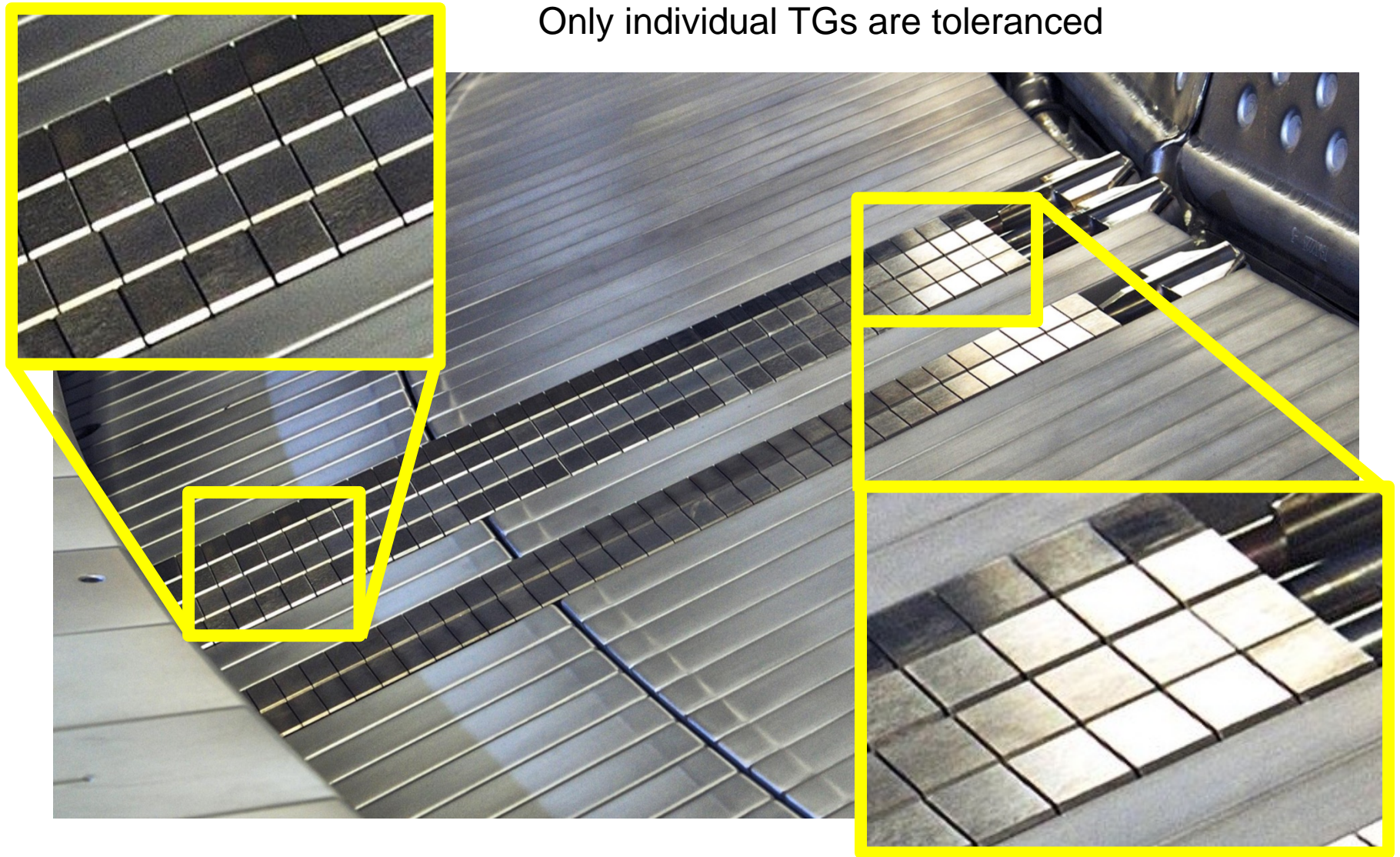


reference 0.5 mm toroidal bevel  
no poloidal bevel  
worst case misalignments  
TG edge and OHS are visible

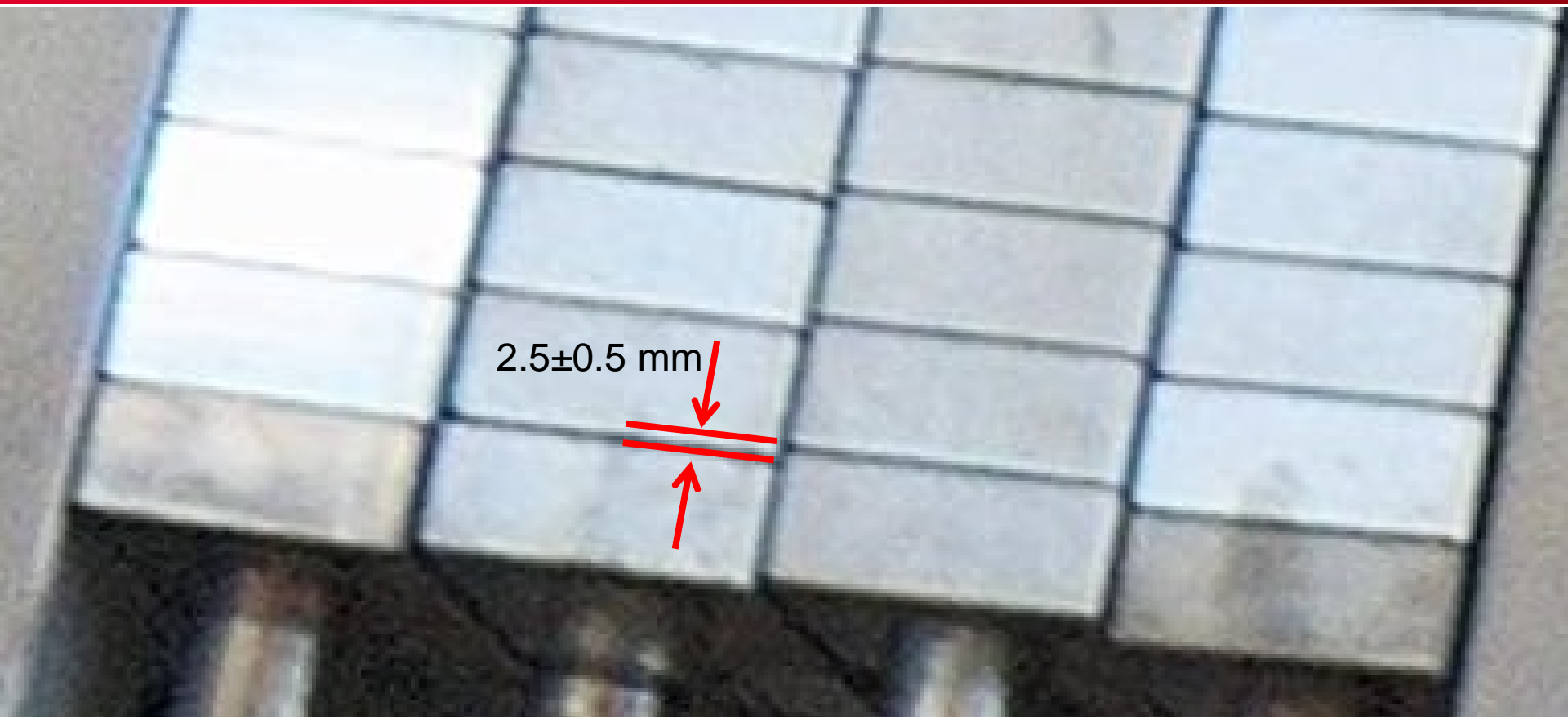
reference 0.5 mm toroidal bevel  
+ additional 0.5 mm poloidal bevel  
→ "shallow poloidal bevel"  
Chosen to shadow TG edge for all possible  
radial misalignments and gap tolerances  
Bonus! → no OHS ... **IF TOROIDAL GAPS  
ARE POLOIDALLY ALIGNED**

# POLOIDAL ALIGNMENT BETWEEN ADJACENT MBS IS NOT SPECIFIED IN WEST (OR ITER) DESIGN

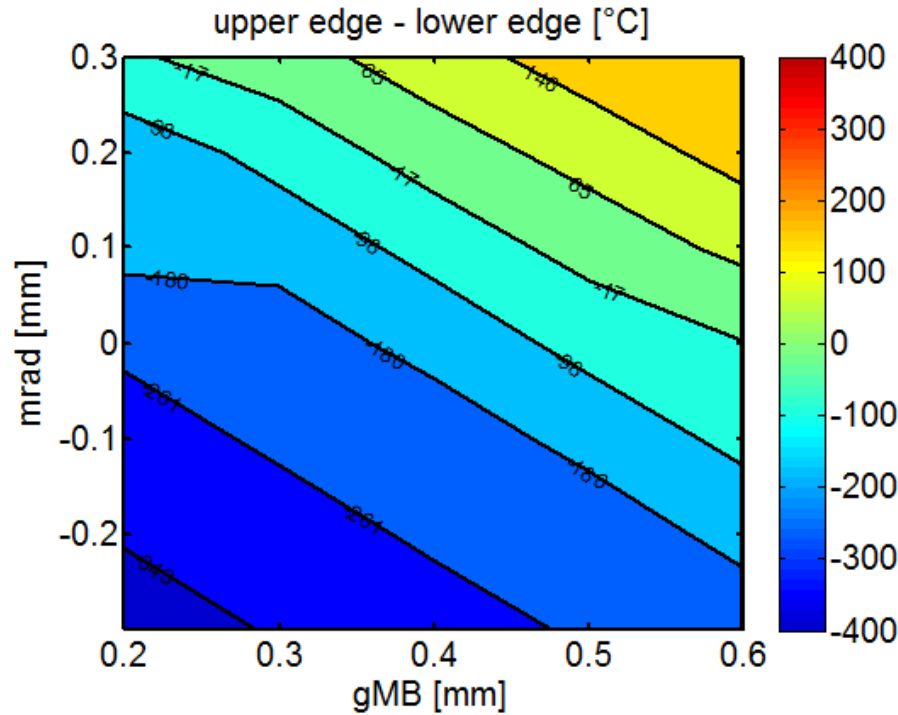
Only individual TGs are toleranced







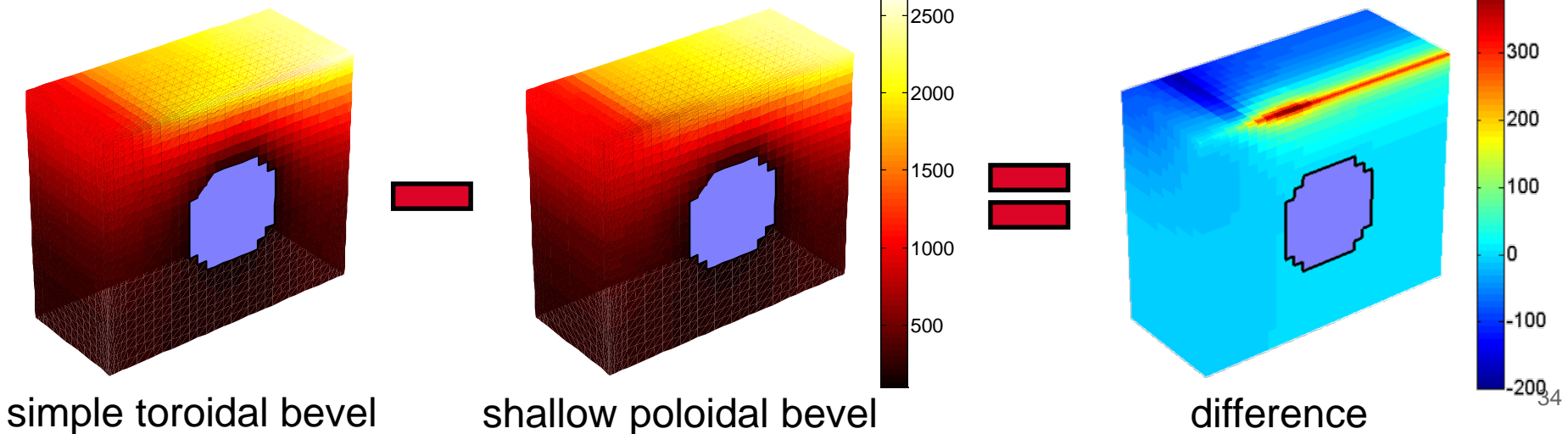
nominal TG width  $g_{MB}=0.5$  mm



negligible increase of top surface heating  
 simple toroidal bevel -  $q_{\text{surf}} / q_{\text{tg}} = 1.56$   
 shallow poloidal bevel -  $q_{\text{surf}} / q_{\text{tg}} = 1.64$

suppression of toroidal edge heating (now cooler than top surface because of shadowing)

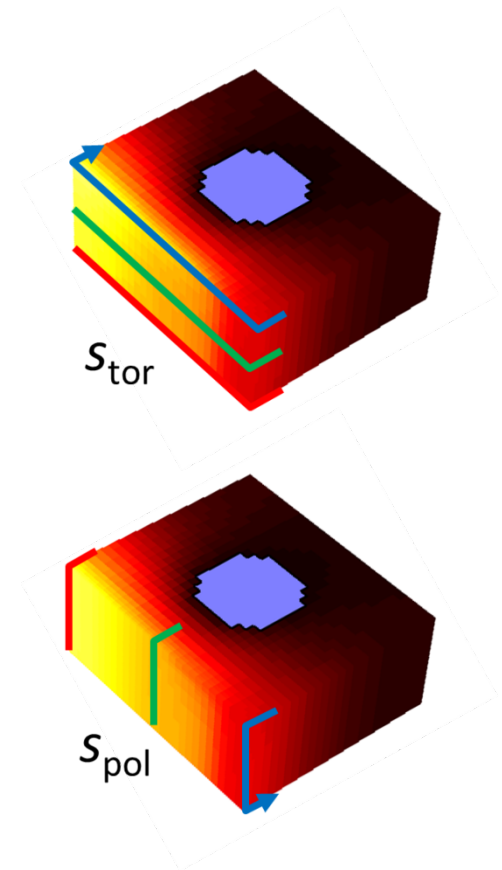
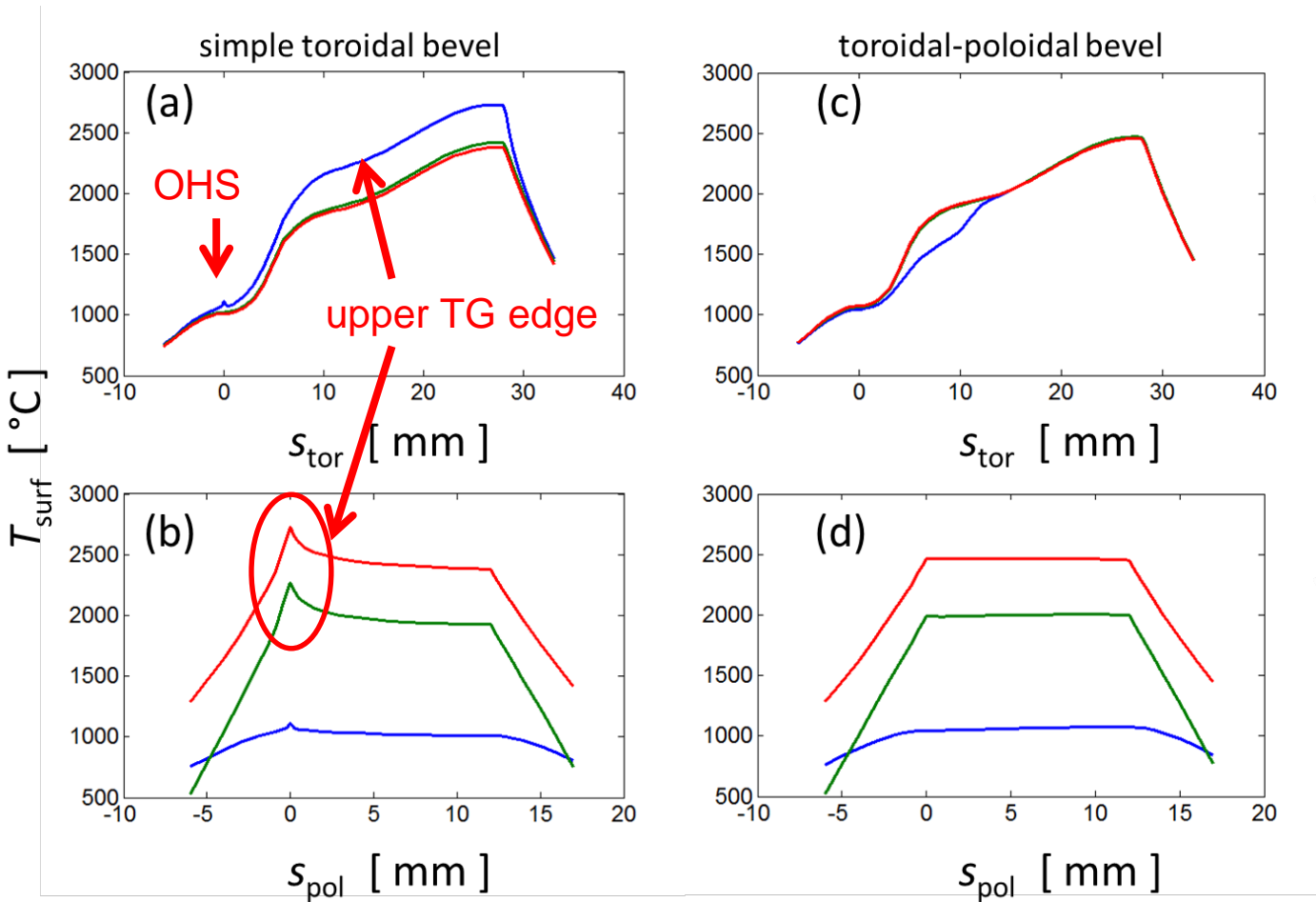
suppression of OHS heating



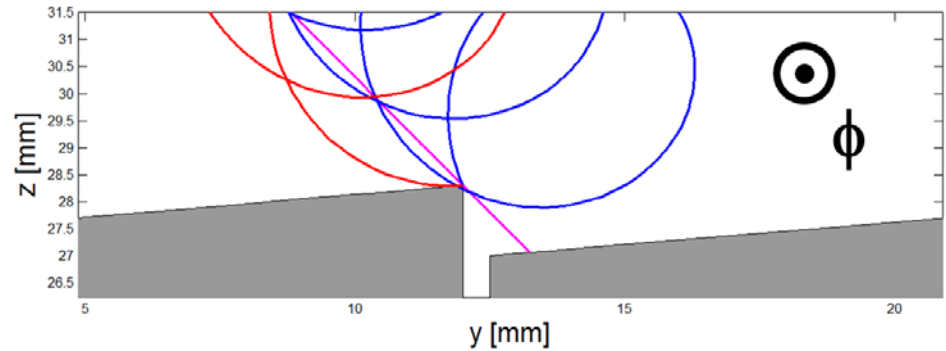
simple toroidal bevel

shallow poloidal bevel

difference

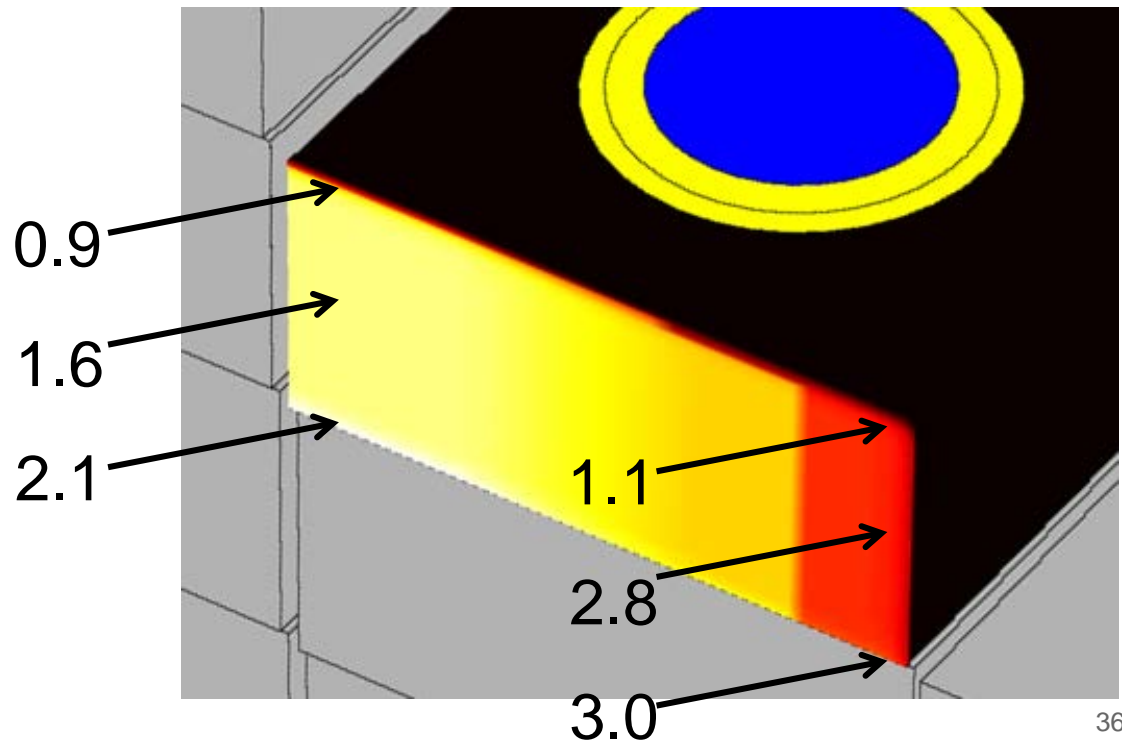
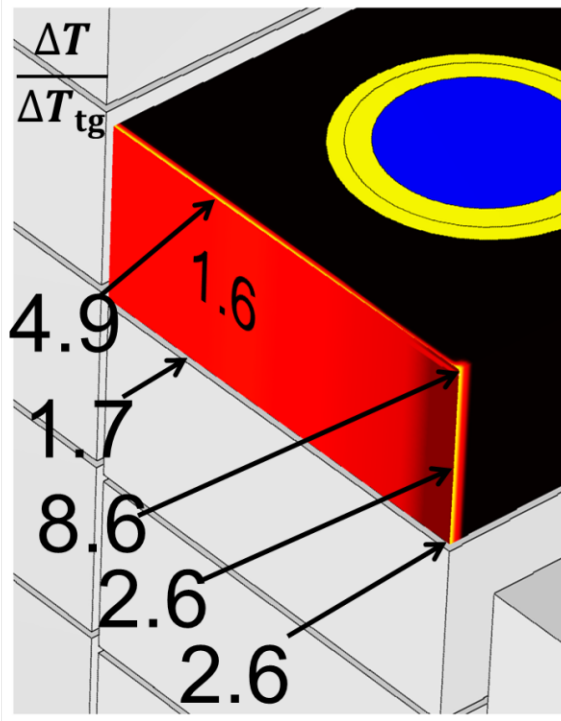


... at the expense of a slight increase along the lower edge



*simple toroidal bevel*

*toroidal-poloidal bevel*



According to ion orbit modelling (and PIC), uncontrolled ELMs will melt all monoblock surfaces and edges at both vertical targets in burning nuclear scenario.

Exposed points ( $<1 \text{ mm}^2$ ) at optical hot spot will be melted or even vapourized.

Edge melting is possible in half-field pre-nuclear scenario.

The reason: a combination of plasma physics (Larmor radius), geometry (enhancement of heating  $\times 2$  at edges,  $\times 3$  at corners), and high MB temperatures.

The simple toroidal bevel solution has been retained for ITER.

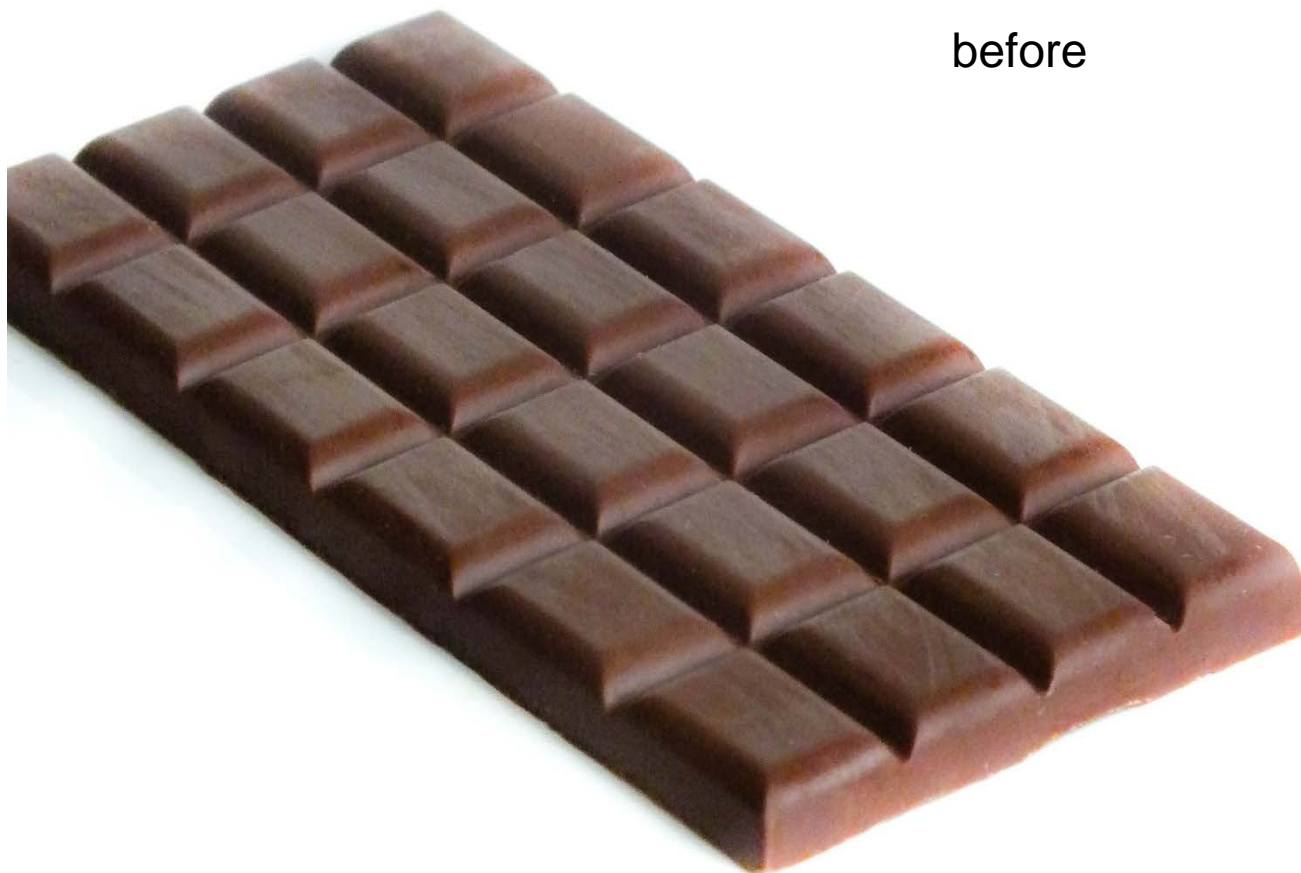
It is too late and would be too expensive to implement a more complex outer target shaping solution at this stage.

In any case there is no solution at the inner target (because of ion Larmor effect)

These findings will be useful for divertor design in future fusion devices.

(Detailed analysis submitted "soon" to Nuclear Fusion journal)

*It is imperative to find ELM-free regimes in ITER.*



before

after

